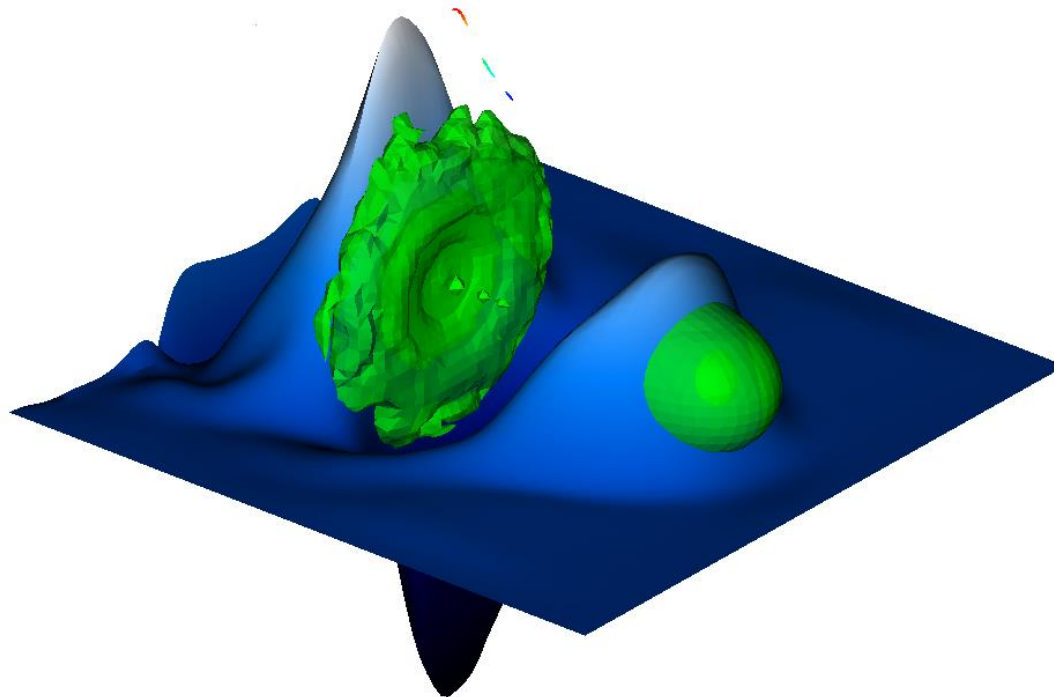


Towards Hybrid Plasma-based Multi-color FELs

Bernhard Hidding^{1,2,3}

¹ Scottish Center for the Application of Plasma Accelerators, University of Strathclyde,

² University of Hamburg & DESY, ³ Department of Physics and Astronomy, UCLA



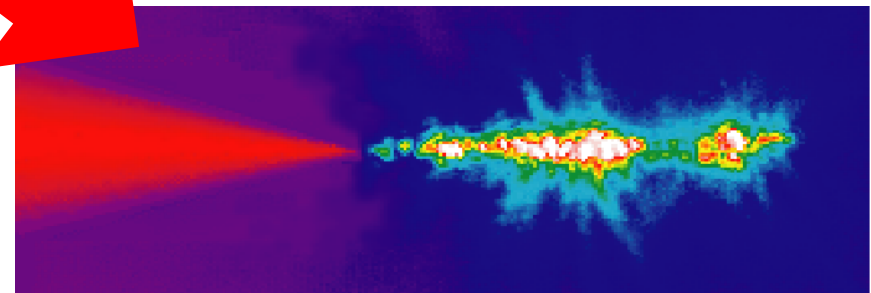
Shrinking accelerators from km to cm: Plasmas



Multiple static metallic cavities
w/ electric fields of ~ 50 MV/m



Single co-propagating plasma cavity
w/ electric fields of ~ 50 GV/m



0 100 200 300
channel length / μm

Prehistoric days: Plasma Wakefield Acceleration

Rutherford/Geiger 1911

World's first particle accelerator experiment:
Matter consists of electrons and ions



Langmuir/Tonks 1928

"We shall use the name *plasma* to describe [a] region containing balanced charges of ions and electrons"

CERN 1956

Future particle accelerators:

Accelerate particles via collective fields by separating electrons and ions in plasmas

Veksler, Budker, Fainberg, *Proc. CERN Symp. High Energy Accelerators*, 1956

Project Matterhorn

Description and computation of nonlinear plasma oscillations

J. Dawson, *Phys. Rev.* 113, 383, 1959

UCLA 1979: **LWFA**

Produce transient charge separation in plasma via Laser Electron Accelerator

Tajima & Dawson, *Phys. Rev. Letters* 43, 1979

CPA 1986

Chirped Pulse Amplification to produce intense enough lasers

Strickland & Mourou, *Optics Comm.* 56, 219, 1986

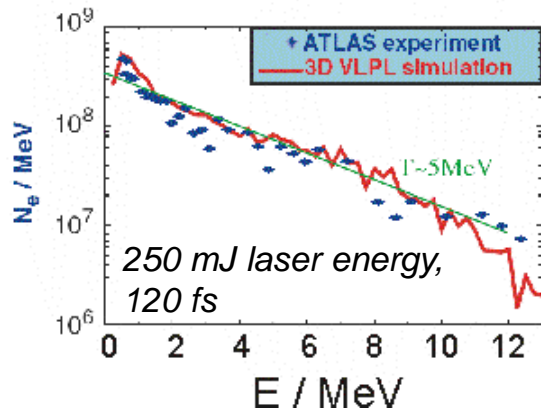
Stanford/UCLA 1985: **PWFA**

Acceleration of Electrons by the Interaction of a Bunched Electron Beam with a Plasma

Chen et al., *Phys. Rev. Letters* 54, 1985

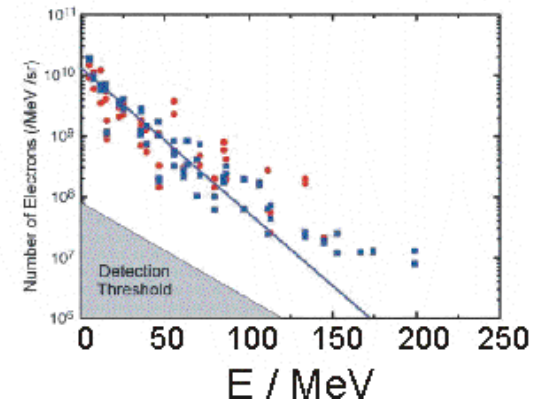
Modern LWFA History

Since 1990s: Exponential beams



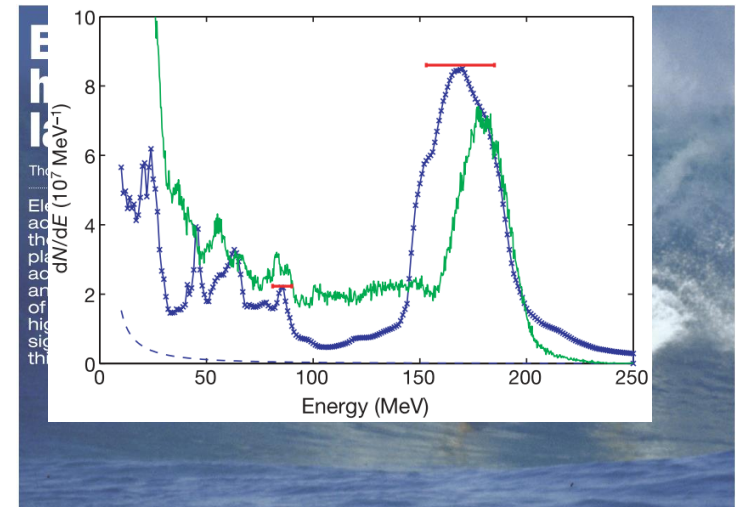
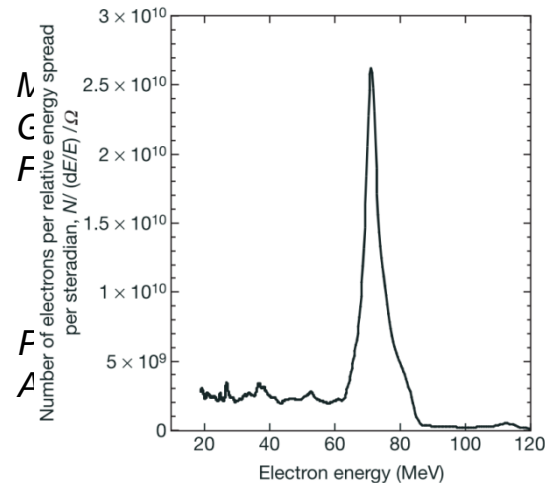
C. Gahn, Phys. Rev. Letters 83, 4772, 1999

1 J laser energy, 30 fs



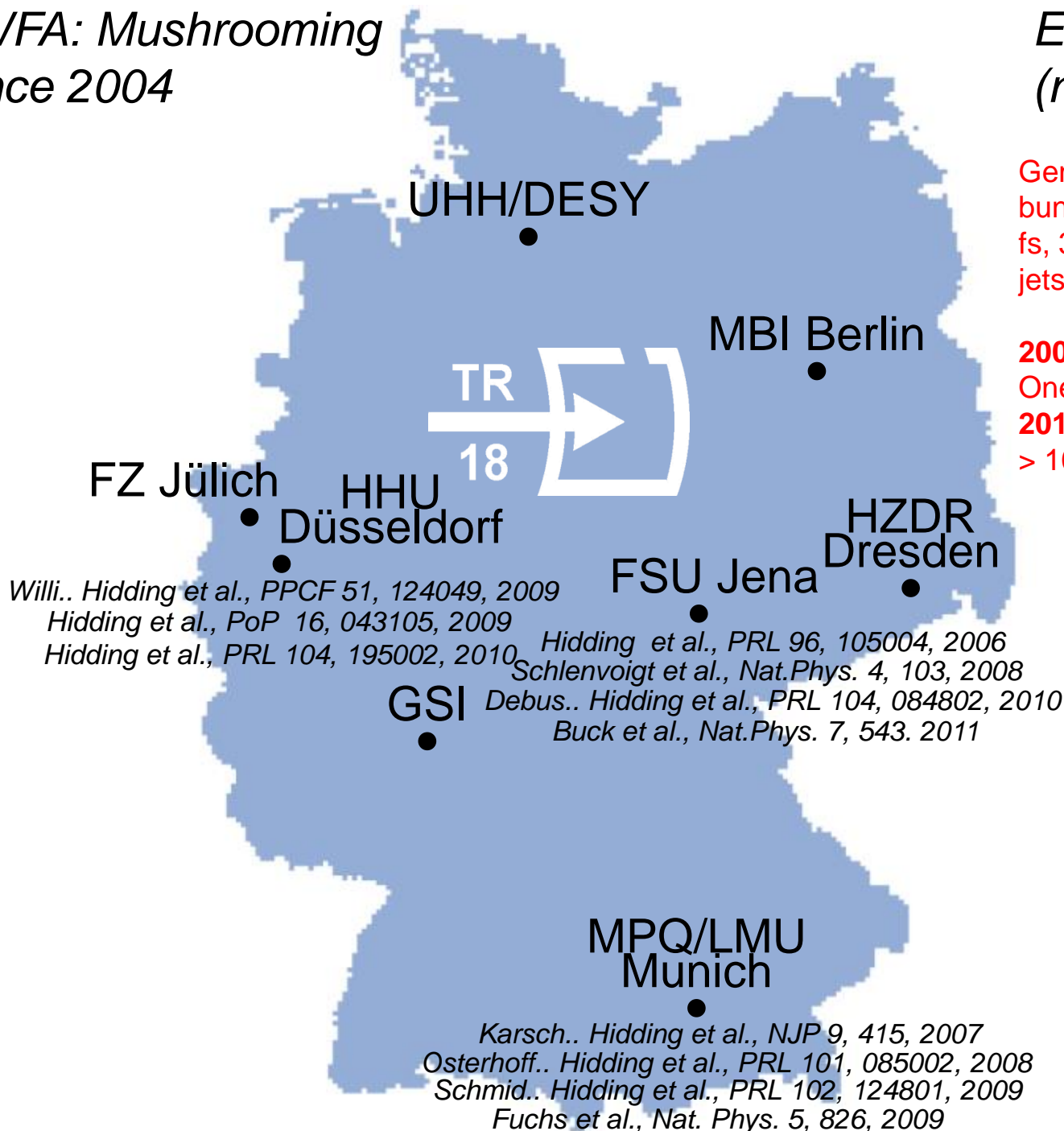
V. Malka, Science 22, 298, 5598, 2002

Since 2004: Quasi-monoenergetic beams



LWFA: Mushrooming since 2004

Example: Germany
(non-exhaustive!)



Generation of μm -scale electron bunches up to 1 GeV with 8-80 fs, 30 mJ-3 J laser pulses in gas jets, capillaries and gas cells

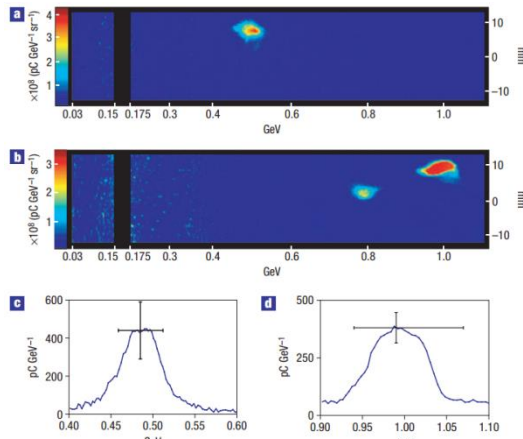
2004:

One laser system with 7 TW

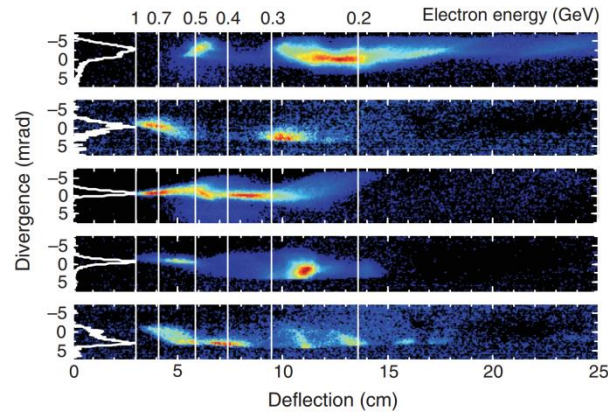
2013:

> 10 laser systems w/ > 100 TW

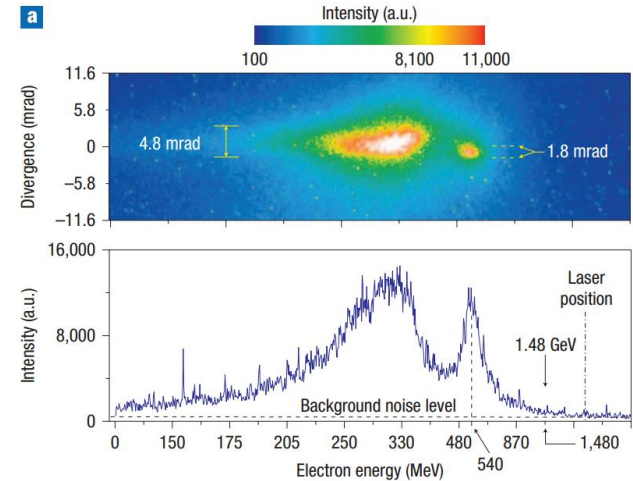
But: Limited beam quality & stability



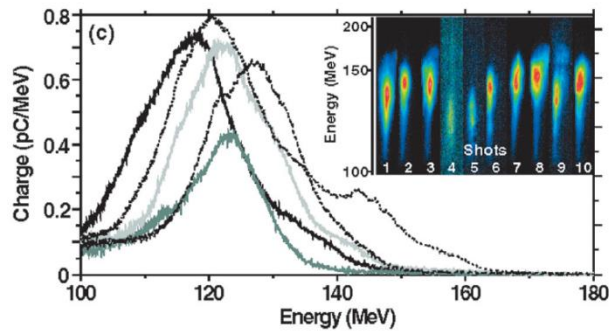
Leemans et al., Nat. Phys. 2006



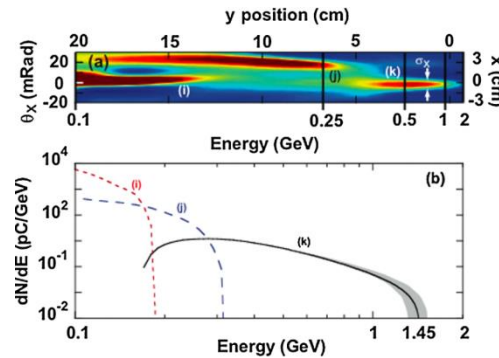
Karsch et al., NJP 9, 415, 2007



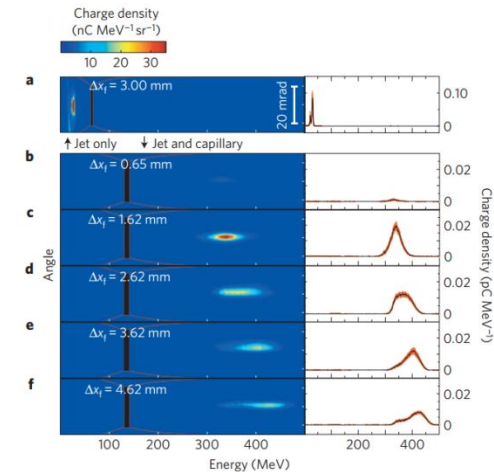
Hafz et al., Nat. Phot. 2008



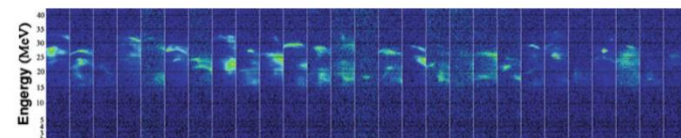
Osterhoff et al., PRL 2008



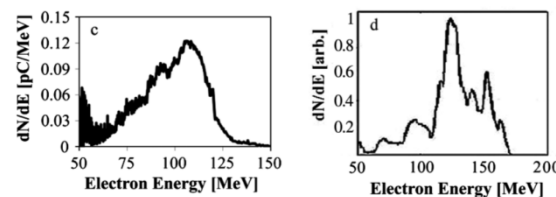
Clayton et al., PRL 2010



Gonsalves et al., Nat. Phys. 2011



Schmid, PhD Thesis 2009



McGuffey et al., PRL 2010

Fundamental Issues of LWFA

Dephasing, Diffraction & Injection limit energy gain and beam quality:

Strategies: Dephasing, Diffraction & Injection

Dephasing:

Use longitudinally tapered plasma profile
Katsouleas, PRA 2056 (1986)

Diffraction:

Use transversally tapered plasma profile
e.g. Hooker et al., JOSA (2000), Leemans et al., Nat. Phys. (2006)

Injection:

External Injection
Clayton et al., PRL 70, 37 (1993) → LAOLA@REGAE (DESY), HZDR, Frascati, France ...

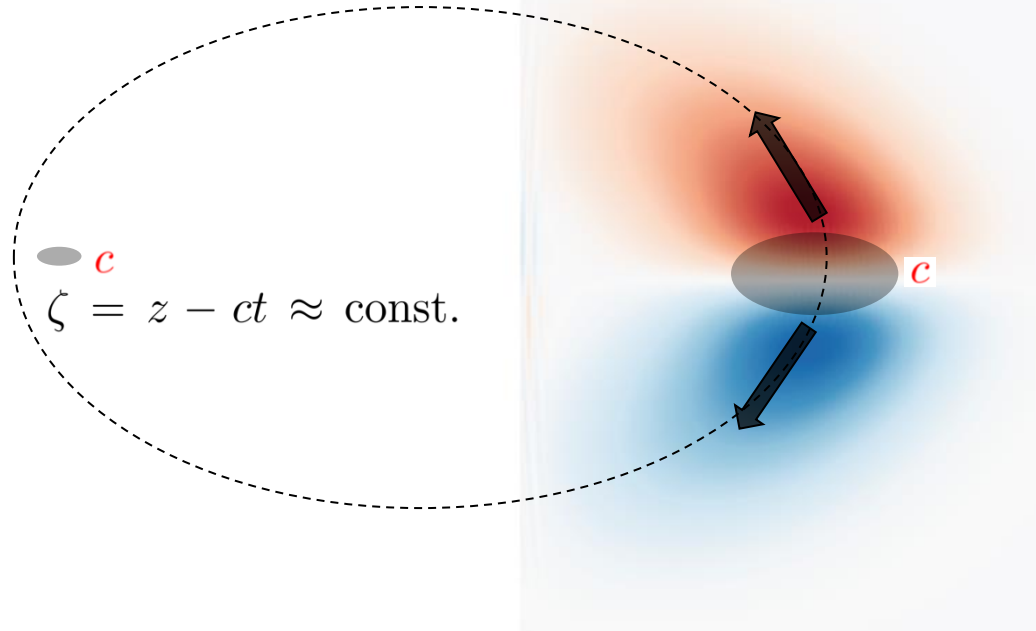
Plasma density transition
Bulanov et al., PRE 58, R5257 (1998), Suk et al., PRL 86, 1011 (2001), Gonsalves et al., Nat. Phys. 7, 862 (2011)

Colliding laser pulses:
Umstadter et al., PRL 76, 2073 (1996), Faure et al., Nature 444, 737 (2006)

Higher state ionization:
Chen et al., JAP 99, 056109 (2006), McGuffey et al., PRL 104, 025004 (2010) (w/ one laser pulse)
Umstadter et al., US Patent 5789876 A (1995), Bourgeois et al., acc. PRL (2013) (w/ two laser pulses)

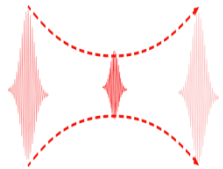
PWFA

„No“ dephasing: (relativistic) driver and accelerated electrons both propagate with $\sim c$
 \Rightarrow witness electrons experience const. (max.) electric field



Much less problems with “diffraction” in PWFA

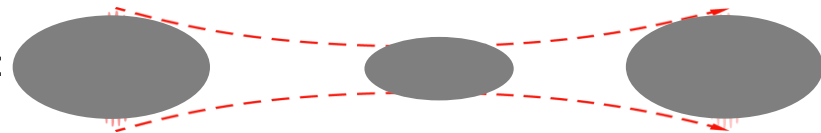
LWFA:



$$\omega(z) = \omega_0 \sqrt{1 + \left(\frac{z}{Z_R} \right)^2} \quad Z_R = \pi \omega_0^2 / \lambda$$

e.g., $Z_R \approx 400 \mu\text{m}$ at $\omega_0 = 10 \mu\text{m}$

PWFA:



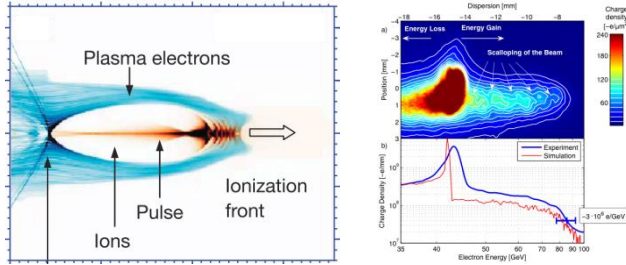
$$\sigma_r(z) = \sigma_{r0} \sqrt{1 + \left(\frac{z}{\beta^*} \right)^2} \quad \beta^* = \sigma_{r0}^2 \gamma / \epsilon_n$$

e.g., $\beta^* \approx 20 \text{ cm}$ at $\sigma_{r0} = 10 \mu\text{m}$,
 $\gamma = 2000$, $\epsilon_n = 10^{-6} \text{ mrad}$

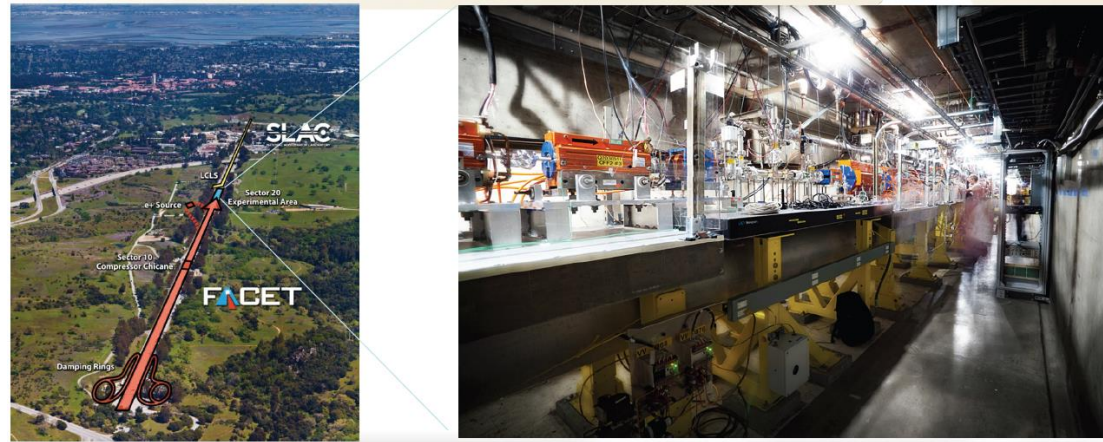


PWFA

SLAC: Energy doubling of 42 GeV electrons in a metre-scale PWFA!



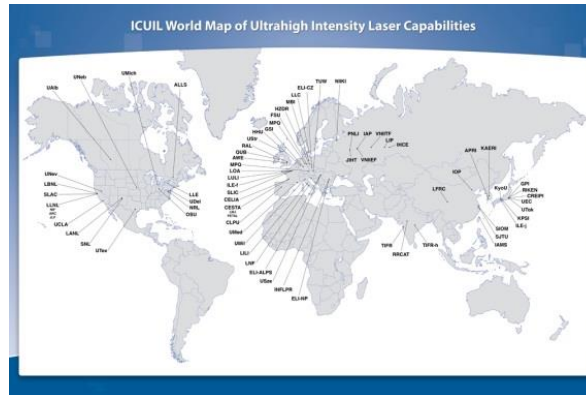
Blumenfeld et al., Nature 445, 741, 2007



But: only one high-energy (> 100 MeV) PWFA facility so far:



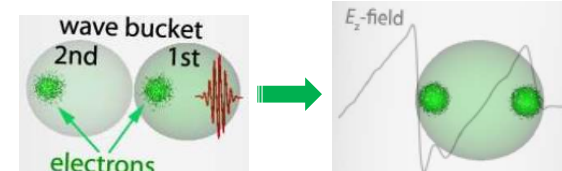
>100 LWFA-capable sites



Idea in 2010:
Use electron bunches from LWFA as particle drivers in subsequent PWFA (afterburner) stage – hybrid LWFA/PWFA

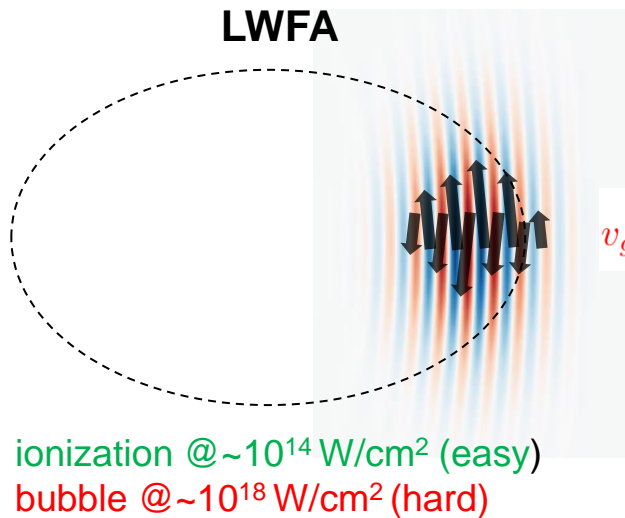
One main reason:
bunches need high current,
need to be compact.

“Monoenergetic Energy Doubling in a Hybrid Laser-Plasma-Accelerator” by using a driver/witness double bunch, Hidding et al., PRL 104, 195002, 2010

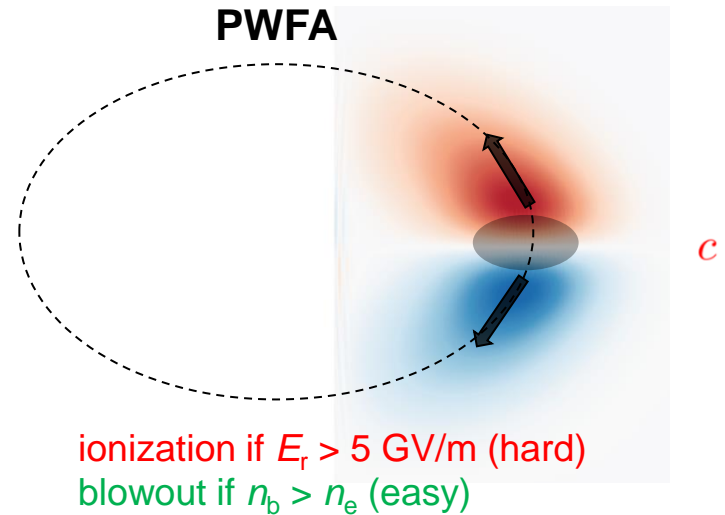


LWFA vs. PWFA

- Laser pulses: transversally oscillating wave, electron bunch: unipolar transverse fields
- Lasers can easily ionize matter, but intensities required to drive a plasma “bubble” orders of magnitude higher
- Electron bunches can drive a plasma “blowout”, but intensities required to self-ionize orders of mag. higher
- PWFA: relativistic electrons move with $\sim c$, no dephasing



$$v_g = c \left(1 - \frac{\omega_p^2}{\omega_0^2} \right)^{1/2} < c$$



- Much longer acceleration distances with relativistic electron bunches (expansion vs. diffraction)

$$\omega(z) = \omega_0 \sqrt{1 + \left(\frac{z}{Z_R} \right)^2} \quad Z_R = \pi \omega_0^2 / \lambda$$

e.g., $Z_R \approx 400 \mu\text{m}$ at $\omega_0 = 10 \mu\text{m}$

free expansion:

$$\sigma_r(z) = \sigma_{r0} \sqrt{1 + \left(\frac{z}{\beta^*} \right)^2} \quad \beta^* = \sigma_{r0}^2 \gamma / \epsilon_n$$

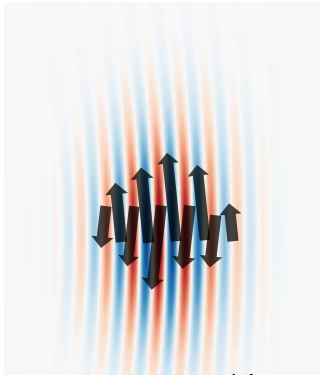
e.g., $\beta^* \approx 20 \text{ cm}$ at $\sigma_{r0} = 10 \mu\text{m}$,
 $\gamma = 2000$, $\epsilon_n = 10^{-6} \text{ mrad}$

Hybrid LWFA & PWFA

Rethink LWFA and PWFA: laser pulses are great for ionization, while electron bunches are better drivers

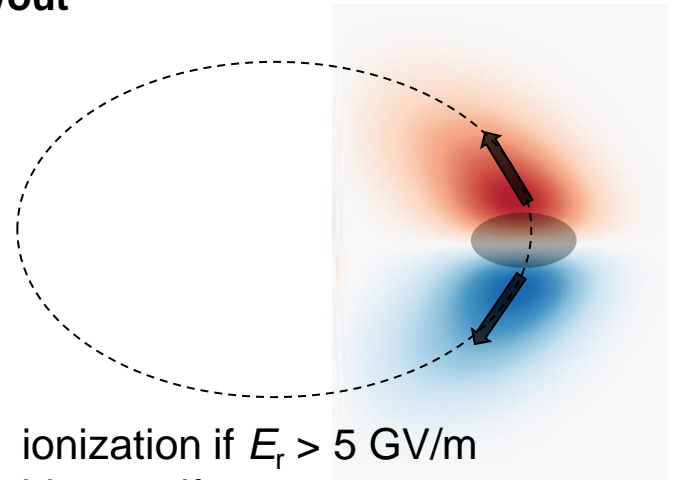
Use the best of both worlds!

Use oscillating fields from laser pulse to ionize and to generate low-transverse momentum electrons



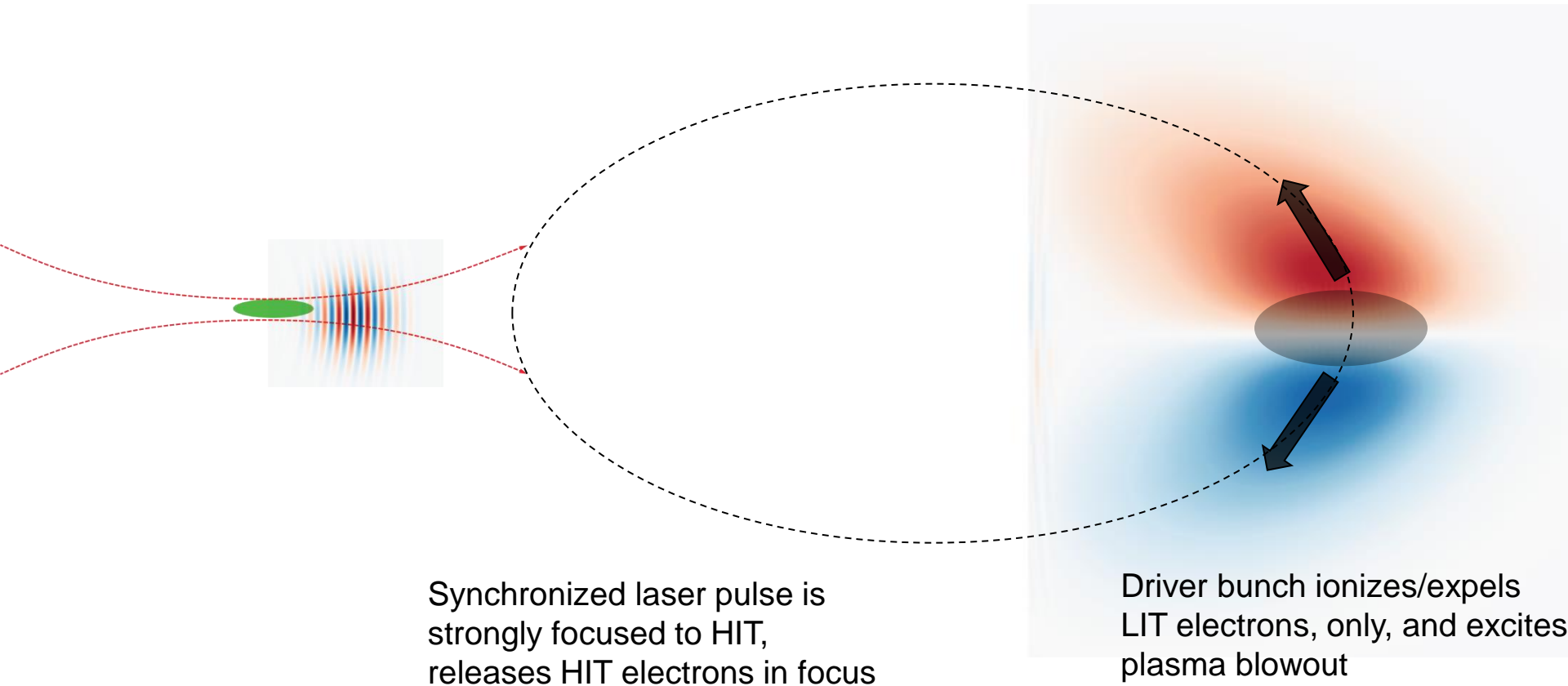
ionization @ $\sim 10^{14} \text{ W/cm}^2$,
produced electrons will receive
very low transverse momentum
(Lawson-Woodward)

Use unidirectional transverse fields from e-bunch to kick out electrons and to excite blowout



ionization if $E_r > 5 \text{ GV/m}$
blowout if $n_b > n_e$

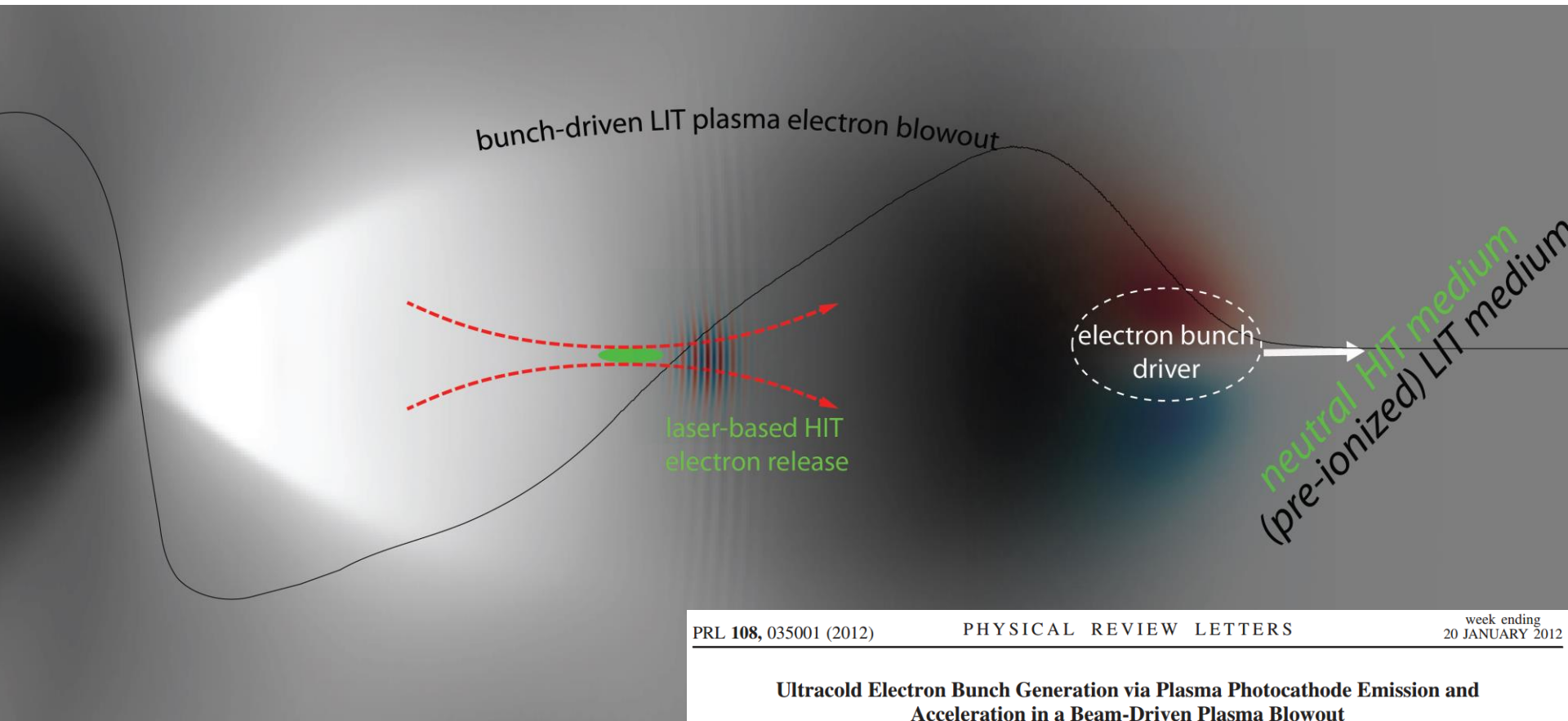
Combine both in media w/ at least two components:
Low-ionization-threshold (LIT), e.g. hydrogen
High-ionization-threshold (HIT), e.g. helium



Beyond Injection: Trojan Horse Plasma Wakefield Acceleration, B. Hidding et al., AAC. APS proc. 2012

Injection: 1590–1600: Latin *injectus* past participle of *in (j) icere* to throw in, equivalent to *in-* + *-jec-* (combining form of *jac-* throw) + *-tus* past participle suffix

Underdense Photocathode PWFA



What's needed:

- LIT/HIT medium
- electron bunch driver to set up LIT blowout
- synchronized, low-intensity laser pulse to release HIT electrons within blowout

Various Potential LIT/**HIT** media candidates

Applicable w/ conventional acc. and LWFA alike

For example

- gaseous H (13.6 eV)/**He (24.6 eV)**
- alkali metals Li, Na, Rb, Cs (~5 eV)/**He (24.6 eV)**
- Rb (4.2 eV)/**Rb⁺ (27.3 eV)**
- Cs (3.9 eV)/**Li (5.4 eV)**

You want the lowest ionization thresholds (to decrease the transverse electron momentum), and a reasonable gap between LIT and HIT medium (ionization corridor)

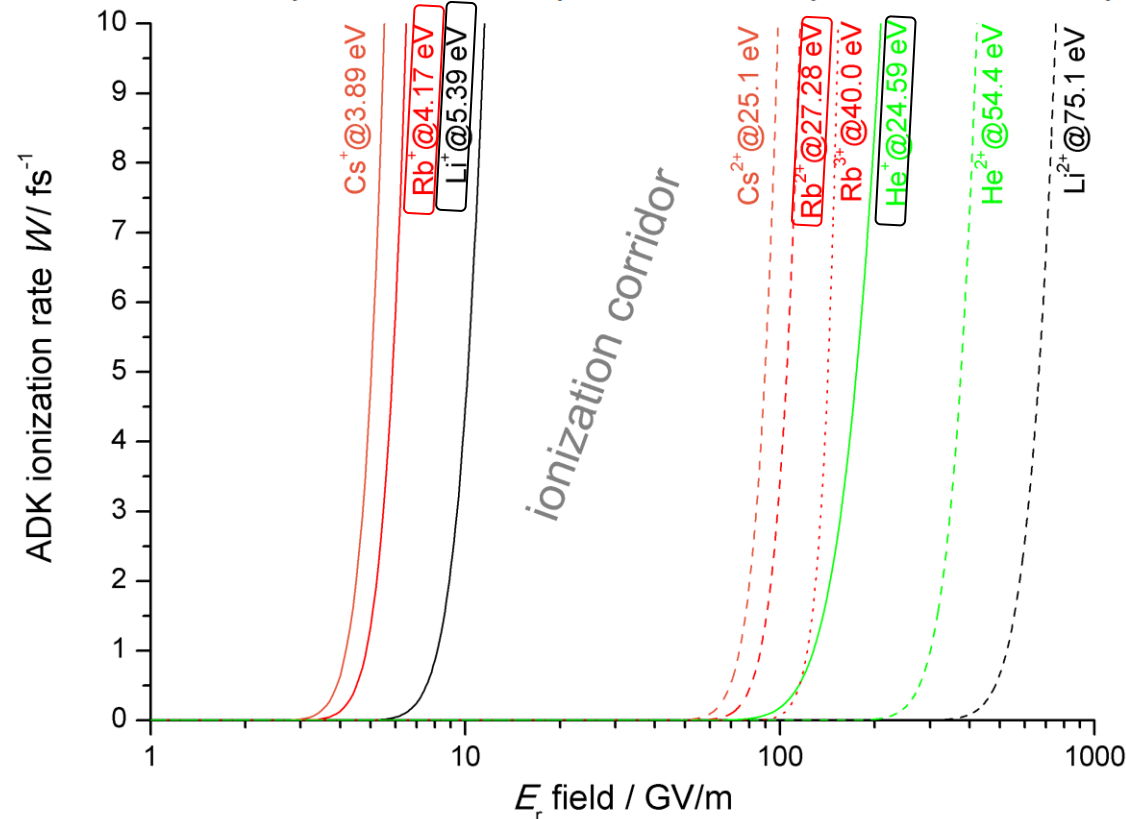
ADK ionization rates: $W(\text{s}^{-1}) \approx 1.52 \times 10^{15} \frac{4^{n^*} \xi_i [\text{eV}]}{n^* \Gamma(2n^*)} \left(\frac{20.5 \xi_i^{3/2} [\text{eV}]}{E [\text{GV/m}]} \right)^{2n^*-1} \times \exp \left(-6.83 \frac{\xi_i^{3/2} [\text{eV}]}{E [\text{GV/m}]} \right)$

Near future @ FACET:

Li (ξ_i = 5.4 eV/He (ξ_i = 24.5 eV)

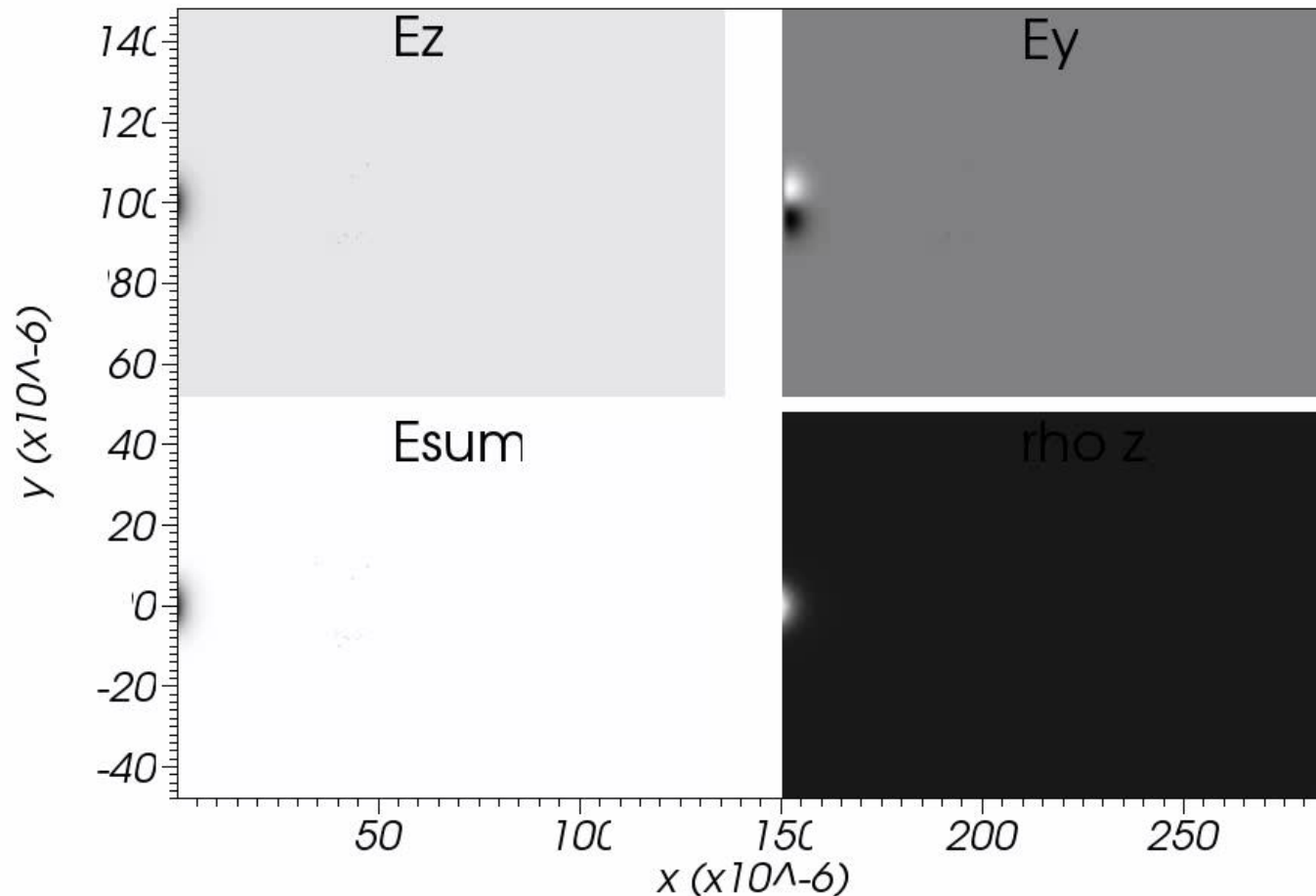
Rb (4.2 eV)/**Rb⁺ (27.3 eV)**

Ar (15.8 eV)



SLAC and FLASH bunches have bunch parameters which are on the verge of self-ionizing alkali metals: R&D in Hamburg on (partial) preionization of alkali metal vapors → plasma lense → assisted self-ionization

Released electrons are compressed, trapped and then co-propagate dephasing-free at the end of the blowout



Both accelerating cavity and photocathode are co-moving in phase with the released electrons

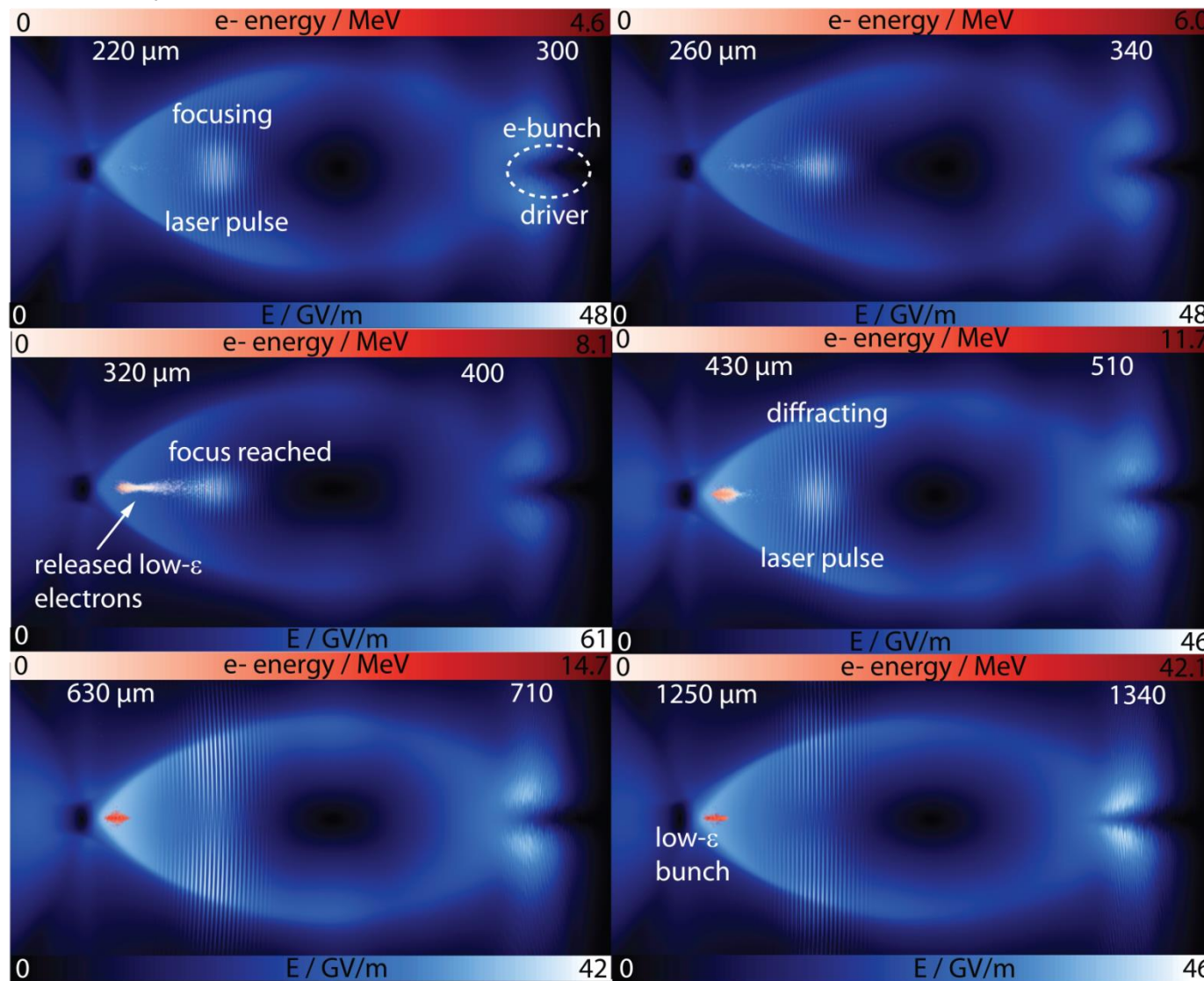
Release laser:
Ti:Sapph, 800 nm,
25 fs, $a_0=0.015$, self-
ionization of LIT medium
by electron bunch

All simulations w/



Laser pulse intensity is crucial

Focus laser pulse intensity has to be just above the ionization threshold of the HIT medium (e.g., helium).



In contrast to LWFA schemes ($\sim 10^{18}$ - 10^{19} W/cm²), here the laser pulse intensity is of the order of $\sim 10^{14}$ - 10^{15} W/cm², **only**.

→ Transverse momentum of bunch electrons is very low → direct consequences for emittance & brightness!

What's the obtainable emittance (in collinear geometry)?

Rough estimation of laser contribution to normalized emittance:

$$\epsilon_n \approx \sigma_{r,HIT} \sigma_{p_r,HIT} / (mc) \approx \boxed{w_0 a_0} / 2^{3/2}$$

ω_0 : laser focus size, a_0 : laser potential

Barrier Suppression Ionization

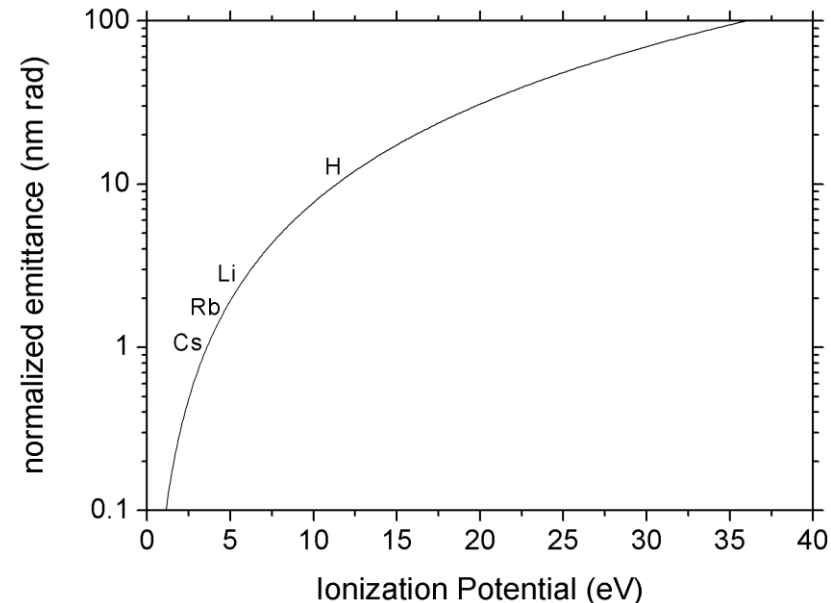
$$I_{BSI} [\text{W/cm}^2] \approx 4 \times 10^9 \frac{\xi_i^4 [\text{eV}]}{Z^2} \quad a_0 = \left(\frac{I}{2\epsilon_0 c} \right)^{1/2} \frac{e\lambda}{\pi m_e c^2}$$

B. Hidding et al., PRL 108, 035001, 2012

Y. Xi et al., PRSTAB 2013

DE patent 2011, US patent 2012

HIT medium	ionization potential	$I_{BSI} / \text{W cm}^{-2}$ @ 800 nm	a_0 at BSI threshold	$\epsilon_{n,x} \approx \omega_0 a_0 / 2.8$ $\omega_0 = 5 \mu\text{m}$
Cs	3.9 eV	9.2×10^{11}	0.00065	$1.2 \times 10^{-9} \text{ m rad}$
Rb	4.2 eV	1.2×10^{12}	0.00075	$1.3 \times 10^{-9} \text{ m rad}$
Li	5.4 eV	3.4×10^{13}	0.00126	$7.1 \times 10^{-9} \text{ m rad}$
H	13.6 eV	1.4×10^{14}	0.008	$1.4 \times 10^{-8} \text{ m rad}$
Cs+	25.1 eV	4.0×10^{14}	0.01362	$2.4 \times 10^{-8} \text{ m rad}$
Rb+	27.3 eV	5.6×10^{14}	0.016	$2.8 \times 10^{-8} \text{ m rad}$
He	24.5 eV	1.4×10^{15}	0.026	$4.6 \times 10^{-8} \text{ m rad}$
Li+	75.6 eV	3.2×10^{16}	0.125	$2.2 \times 10^{-7} \text{ m rad}$



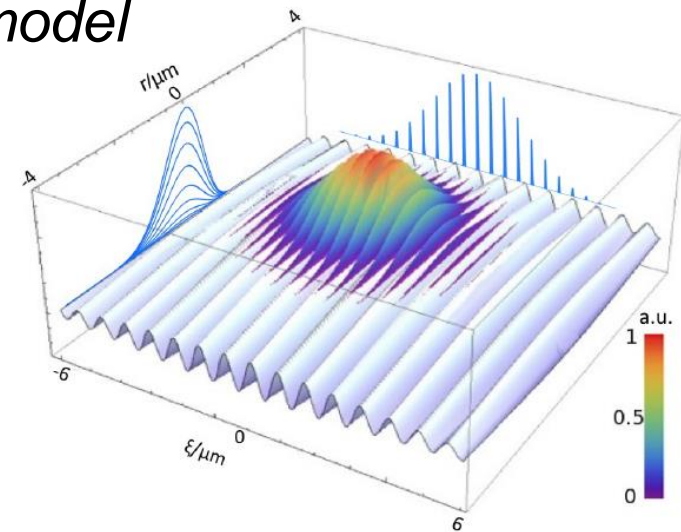
Note:

- Barrier Suppression Ionization is an upper limit
- This is all in laser polarization plane $\rightarrow \epsilon_n$ further decreased in perpendicular plane
- At 800 nm $\rightarrow \epsilon_n$ further decreased w/ higher frequencies $\rightarrow \epsilon_n$ **down to 10^{-10} m rad possible?**

PIC simulation results d'accord with hybrid model

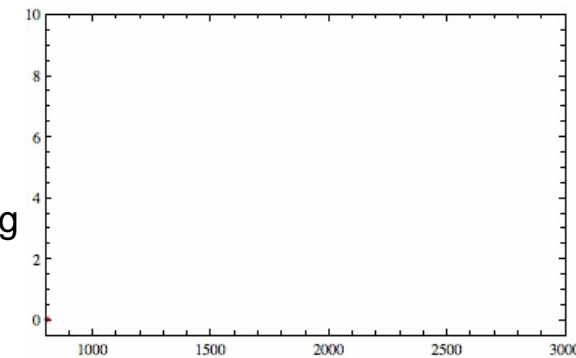
Y. Xi et al., PRSTAB, 2013

- Ionization based on ADK and YI (Yudin-Ivanov-model). G. L. Yudin and M. Y. Ivanov, *Phys. Rev. A*, 64:013409, 2001.
- Detailed numero-analytical analysis shows that $\varepsilon_{n,y}$ is about an order of magnitude lower, and increases slower than $\varepsilon_{n,x}$ as intensity increases. $\varepsilon_{n,y}$ down to the $\varepsilon_{n,y} \approx 10^{-9}$ m rad level or less.



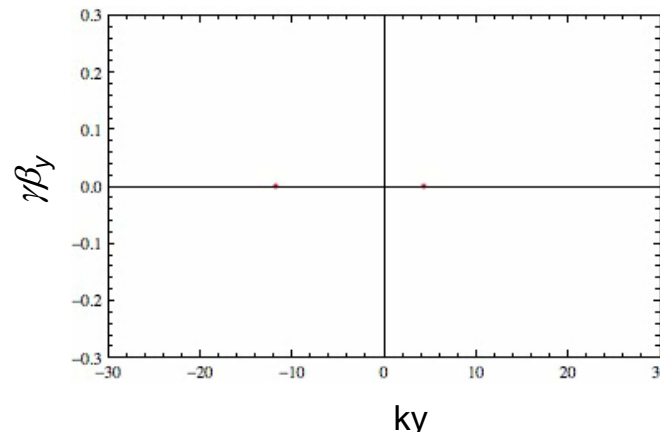
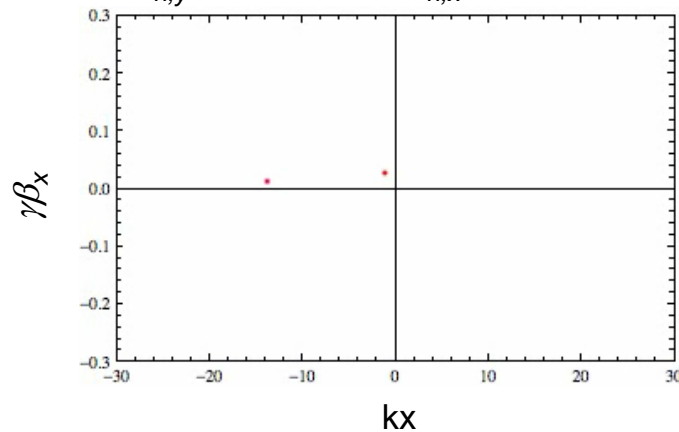
	RF photoinjector	underdense photocathode
beam emittance sources	RF field	ponderomotive motion
	thermal effects	phase mixing
	space charge	space charge

plasma velocity bunching



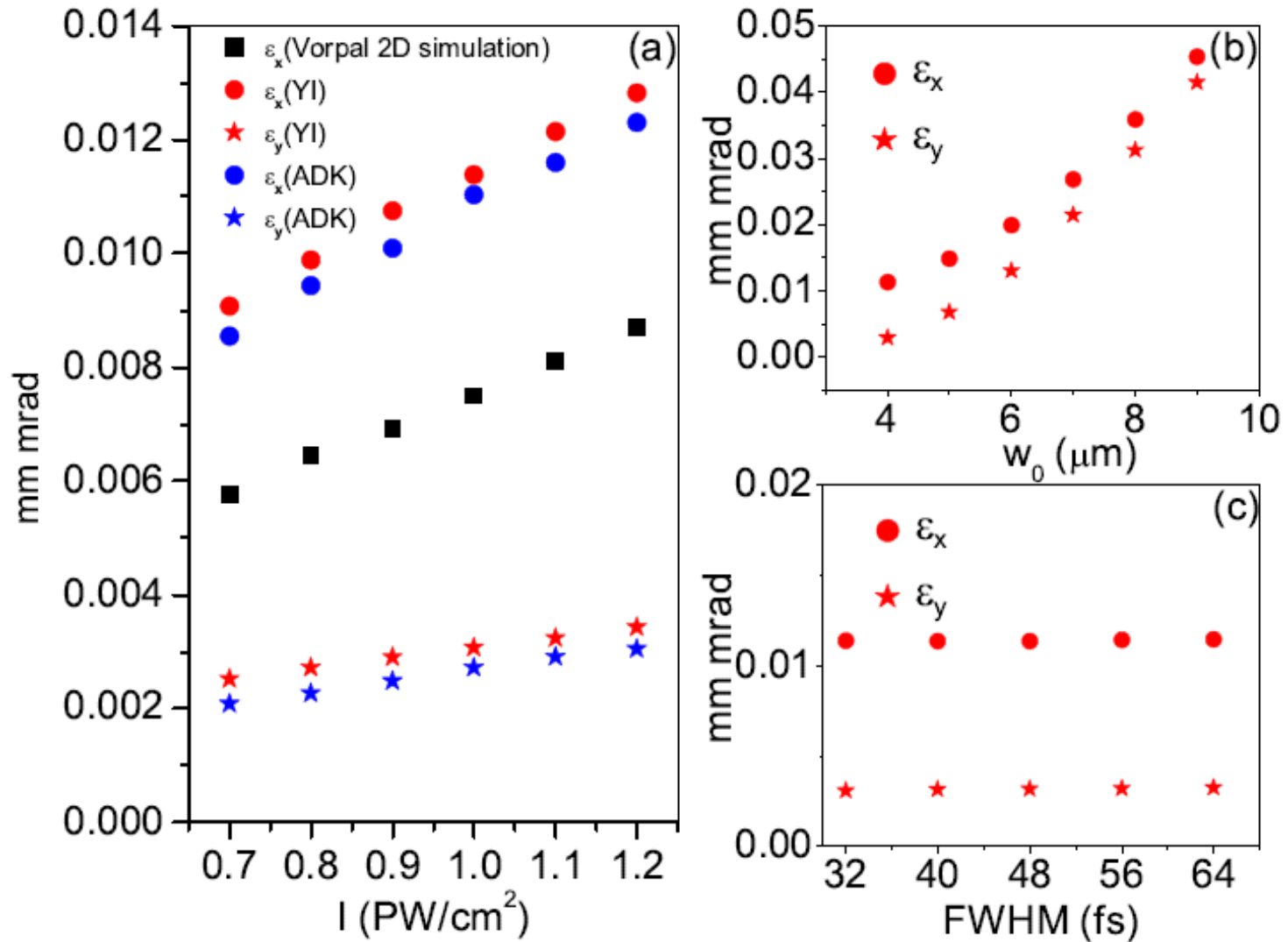
Laser linearly polarized in x-direction:

$\varepsilon_{n,y}$ smaller than $\varepsilon_{n,x}$ due to absence of ponderomotive motion



PIC simulation results d'accord with hybrid model

Y. Xi et al., PRSTAB, 2013



Very good agreement w/
crude laser contrib- only
scaling

$$\epsilon_n \approx \sigma_{r\text{HIT}} \sigma_{p_r, \text{HIT}} / (mc) \approx \boxed{w_0 a_0} / 2^{3/2}$$

Emittance is the key! for FEL, HEP...

- Ultralow emittance yields ultrahigh electron beam brightness even at low charge

e.g., at $Q \approx 1 \text{ pC}$, $\sigma_z \approx 2.5 \text{ } \mu\text{m}$, $I_p \approx 120 \text{ A}$
 $\implies B \approx 2I_p/\epsilon_n^2 \approx 10^{19} - 10^{20} \text{ Am}^{-2}\text{rad}^{-2}$

→ photon brightness

exceeds those of LCLS by wide margin

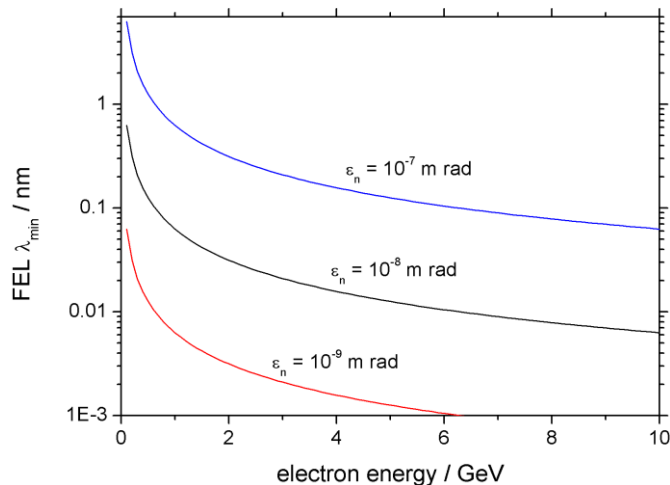
- gain length $L_{g,1D} = \frac{\lambda_u}{4\pi\sqrt{3}\rho_{1D}} \propto B_e^{-1/3}$

exceeds those of LCLS by wide margin

- minimum theoretical FEL wavelength

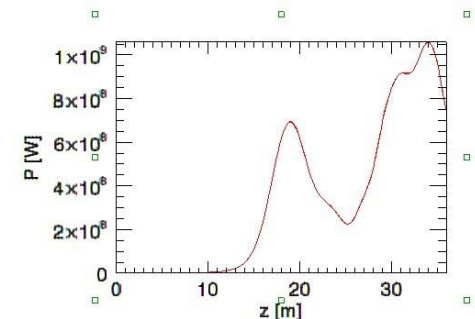
e.g., with $\epsilon_n = 1 \times 10^{-9} \text{ m rad}$ at $1 \text{ GeV} \implies \lambda_{\min} \approx 4\pi\epsilon_n/\gamma < 0.1 \text{ } \text{\AA}$

exceeds those of LCLS by wide margin



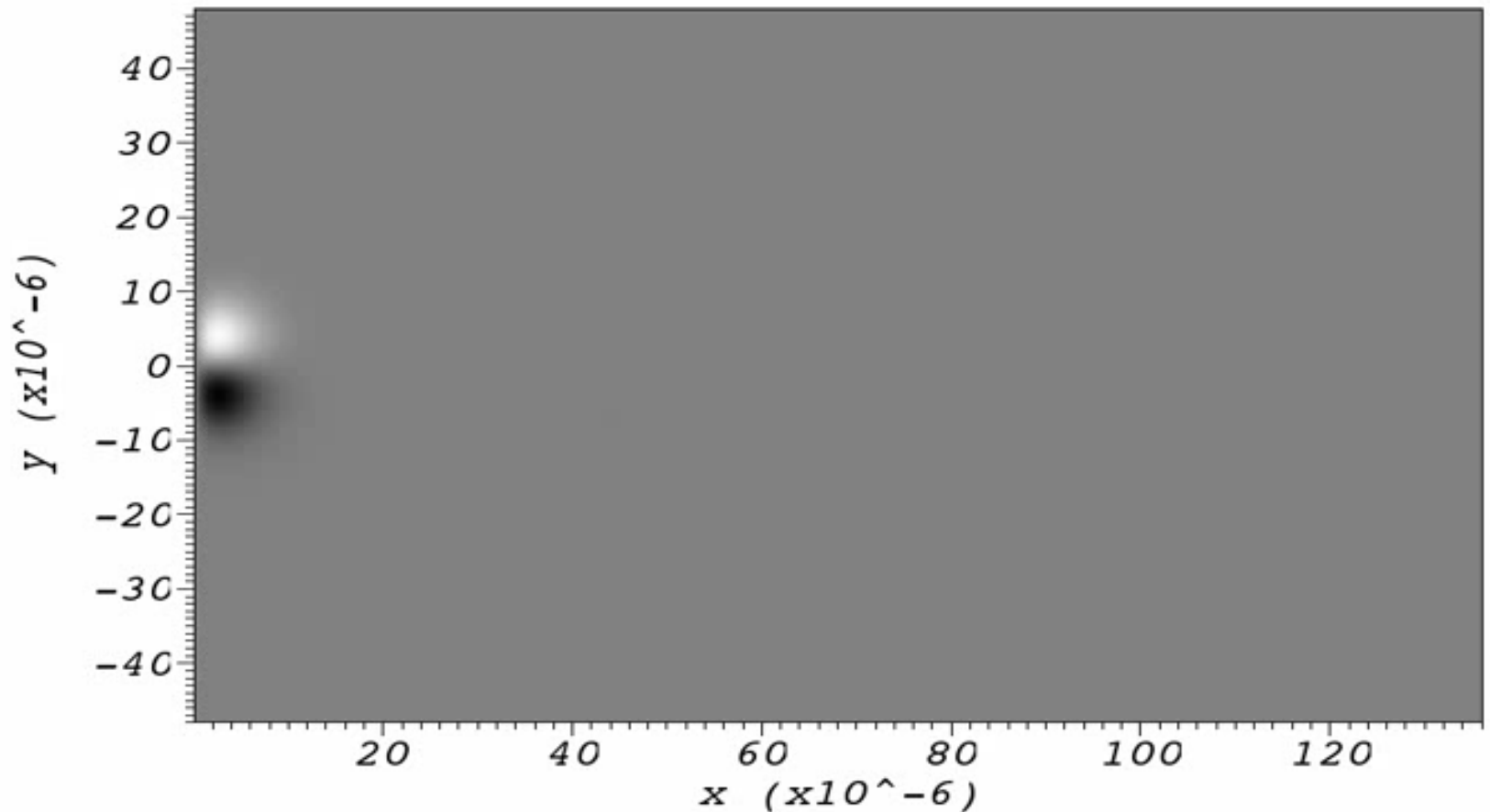
GENESIS calculation for 2 pC bunch, $\epsilon_n = 3 \times 10^{-8} \text{ m rad}$ (only)
w/ “Finndulator” as in O’Shea et al., PRSTAB 13, 070702 (2010),
4.3 GeV: LCLS performance after 20 m!

(First FEL calculations for the Trojan
scenario, done in 2011)



- not “only” for X-ray FEL, also colliders need extremely high beam quality $L = \frac{fN^2}{4\pi\sigma_x\sigma_y}$

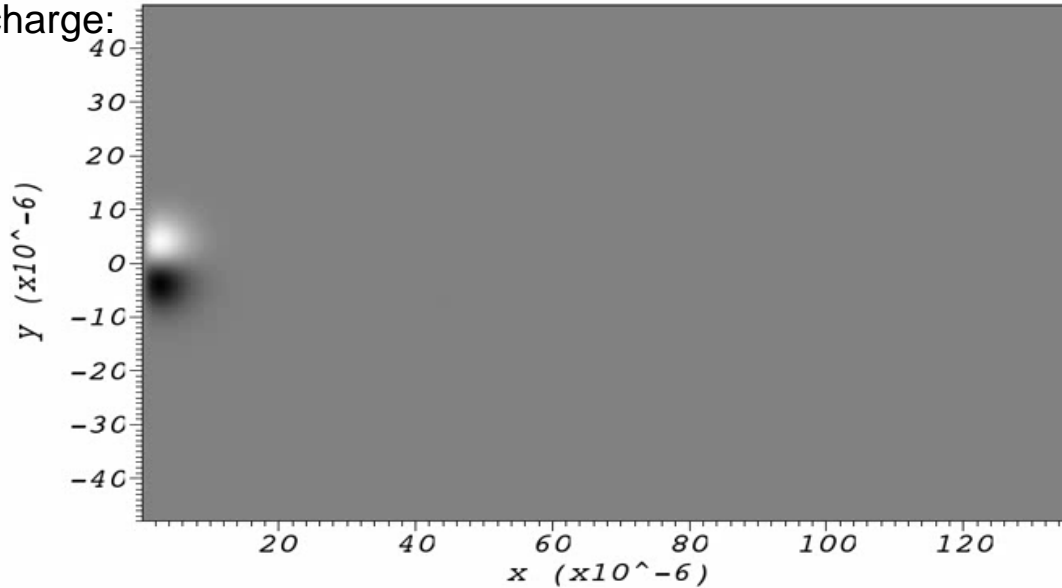
Work at higher laser frequencies, lower intensities and tight focusing: Attosecond electron bunches possible



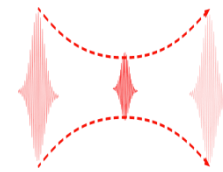
Tunability: charge

Tune charge via laser intensity, spot size and focusing (Rayleigh length)

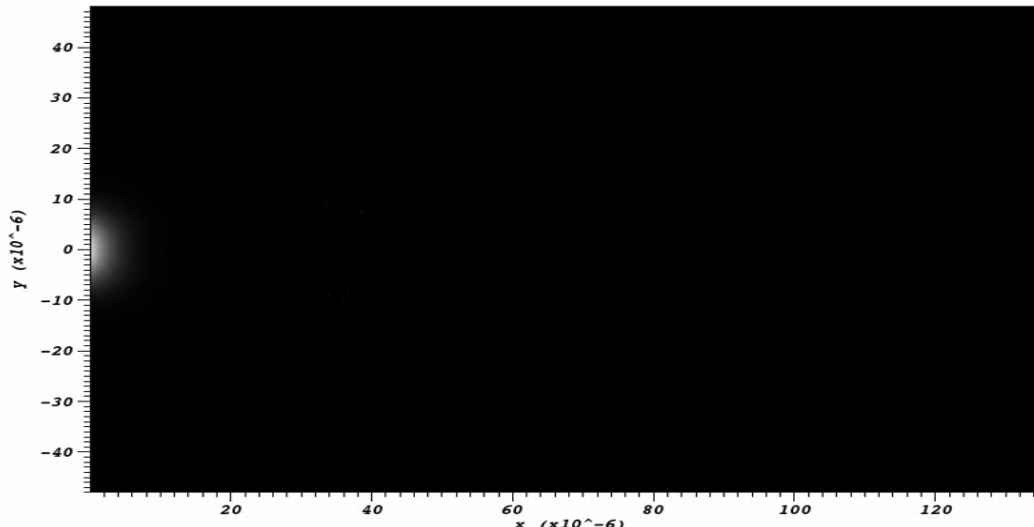
Low charge:



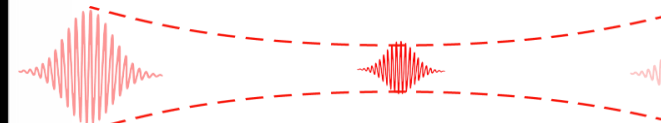
Low laser
intensity, short
Rayleigh length

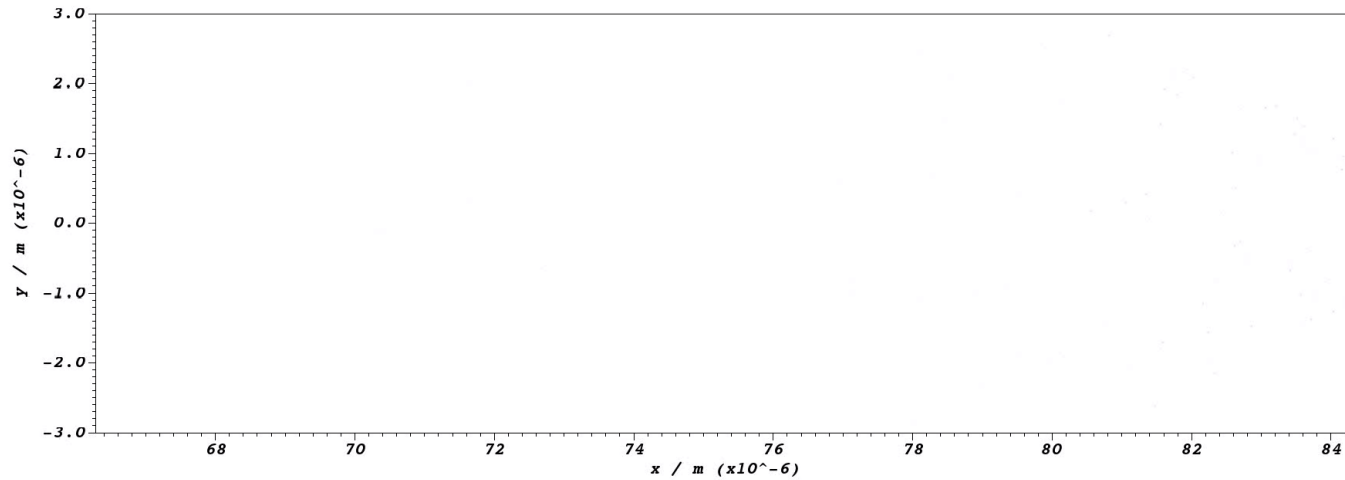


High charge:

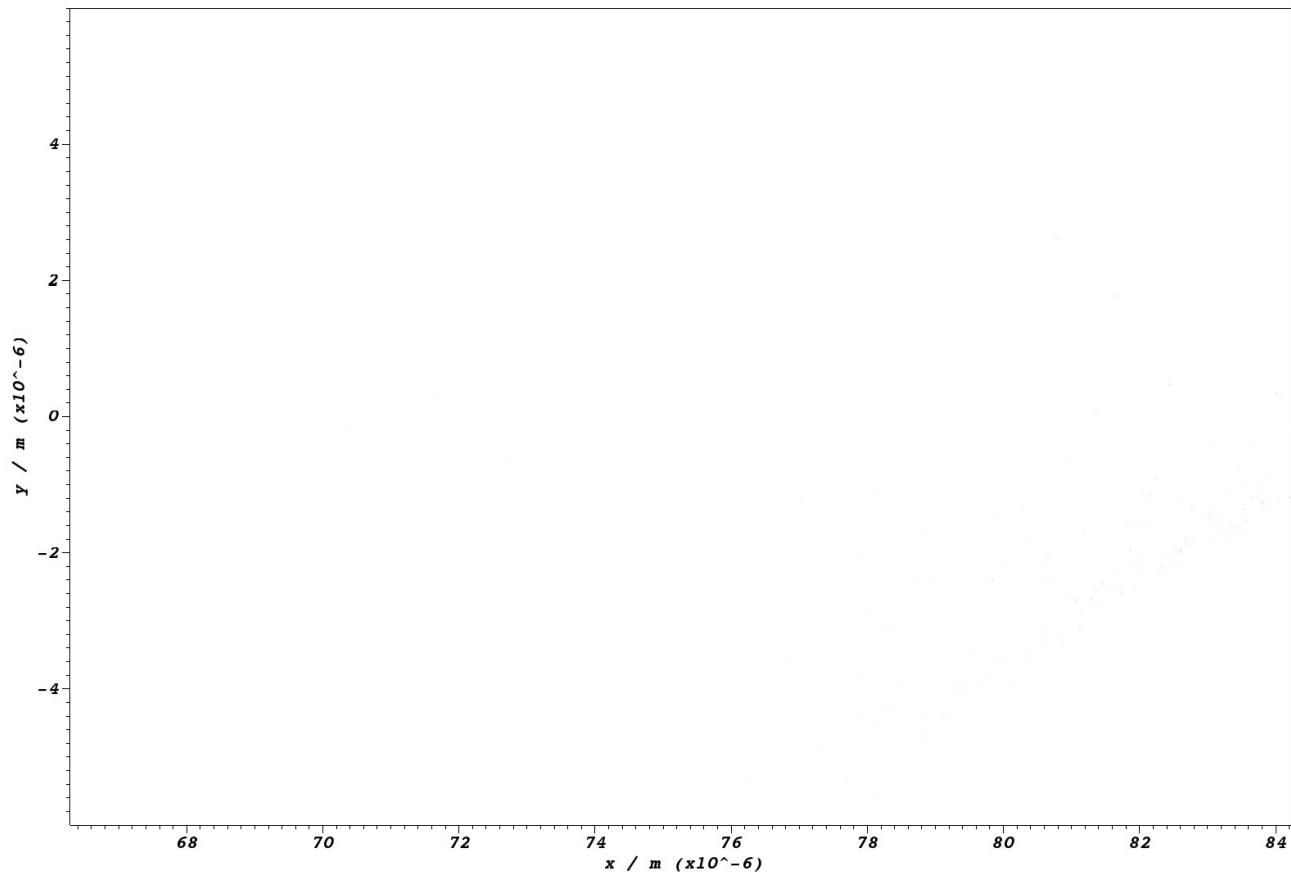


Higher laser
intensity, longer
Rayleigh length





Unique electron fine “snake” structure reflects compressed laser pulse field maxima
 → Coherent radiation + HHG, e.g. in undulator (à la Stupakov et al., PRSTAB 16, 010702, 2013)



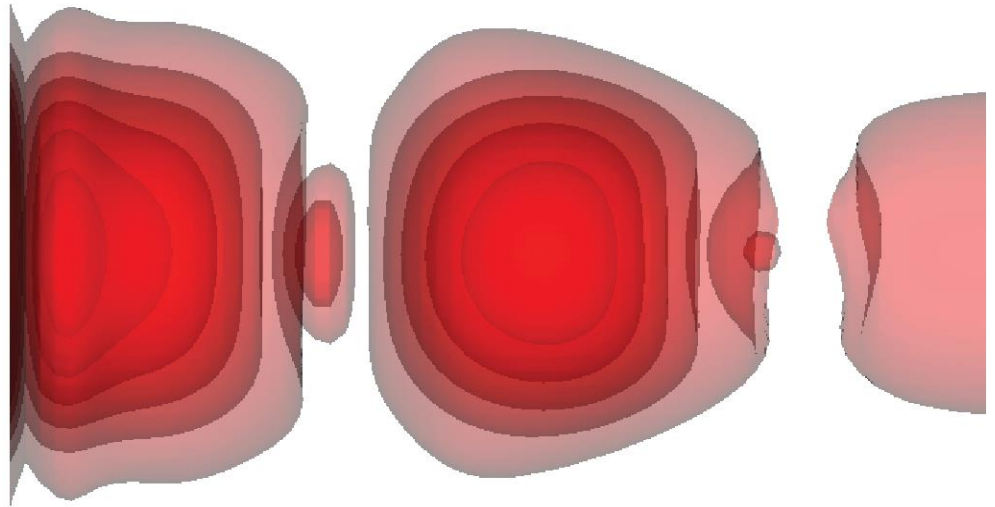
Off axis (e.g., 3 micron) electron release yields controlled betatron oscillations

Multi-bunch production

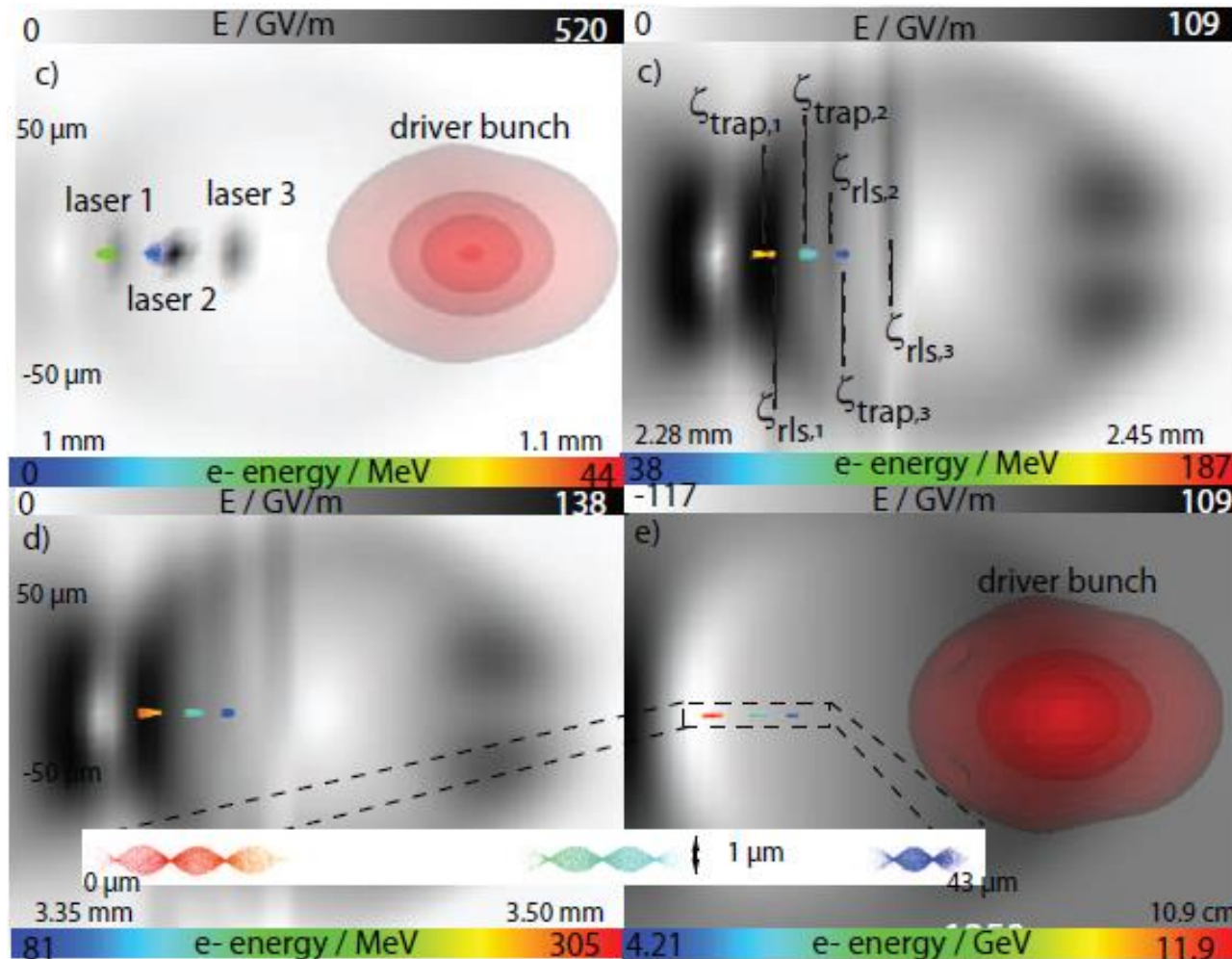
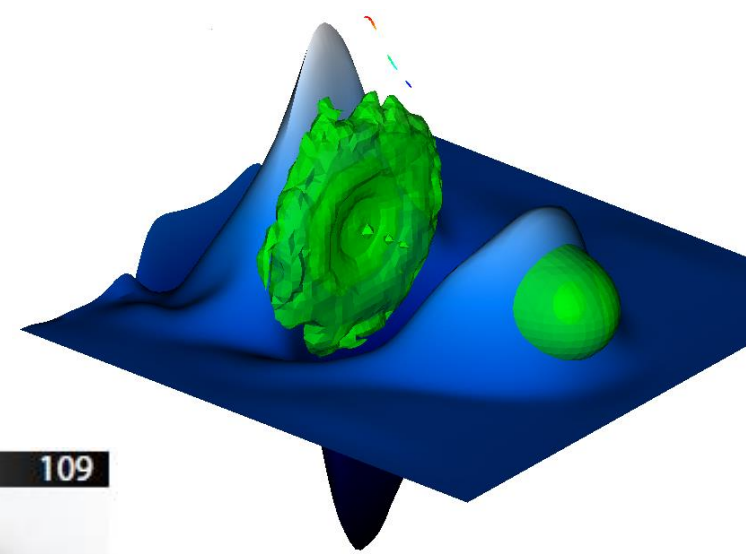
If the wake's potential is large enough:
trapping possible at various positions

Using multiple rfs laser pulses leads
to multi-bunches

FACET parameters, using Plasma/Li+
as LIT/HIT media



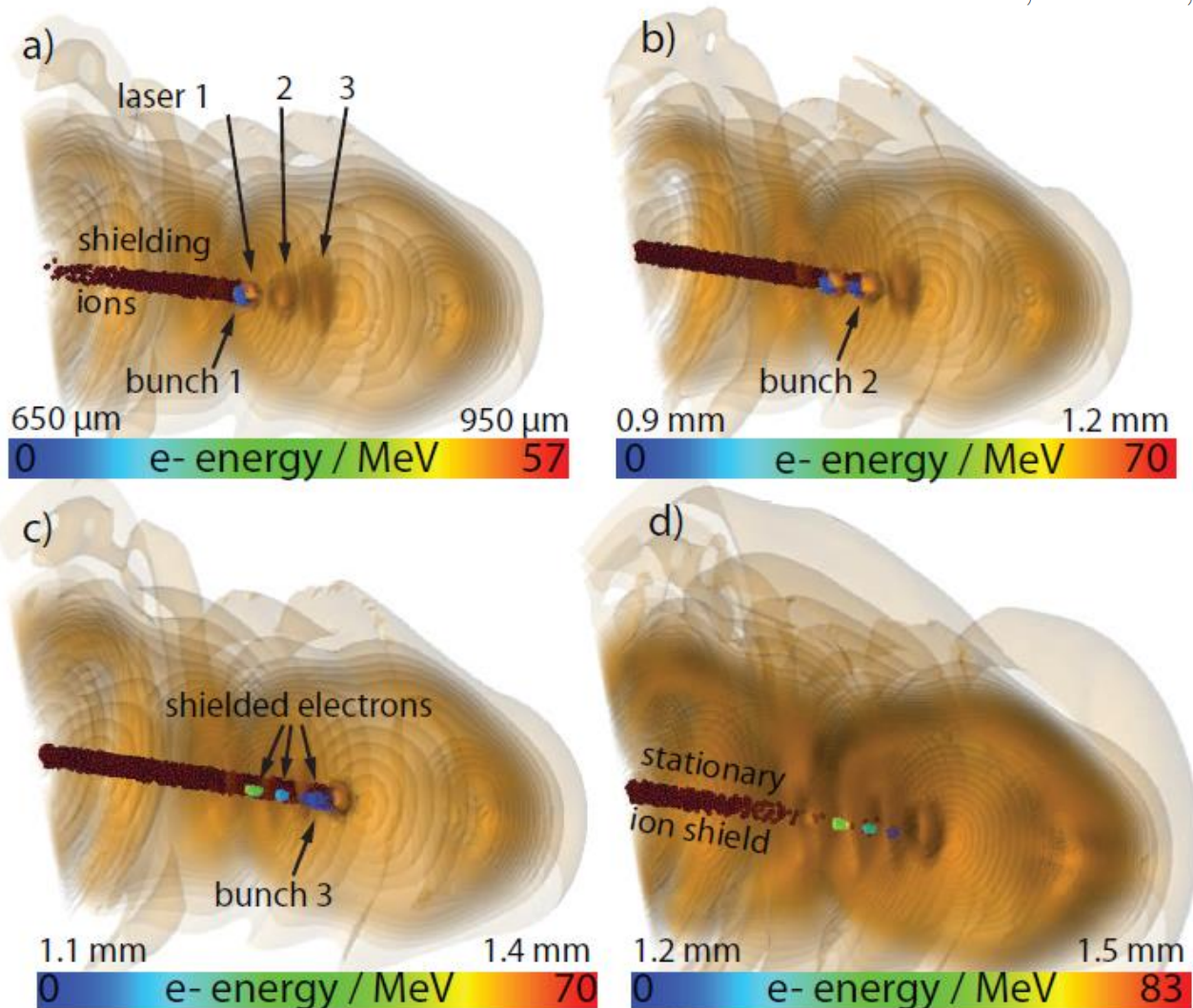
Multi-bunch production



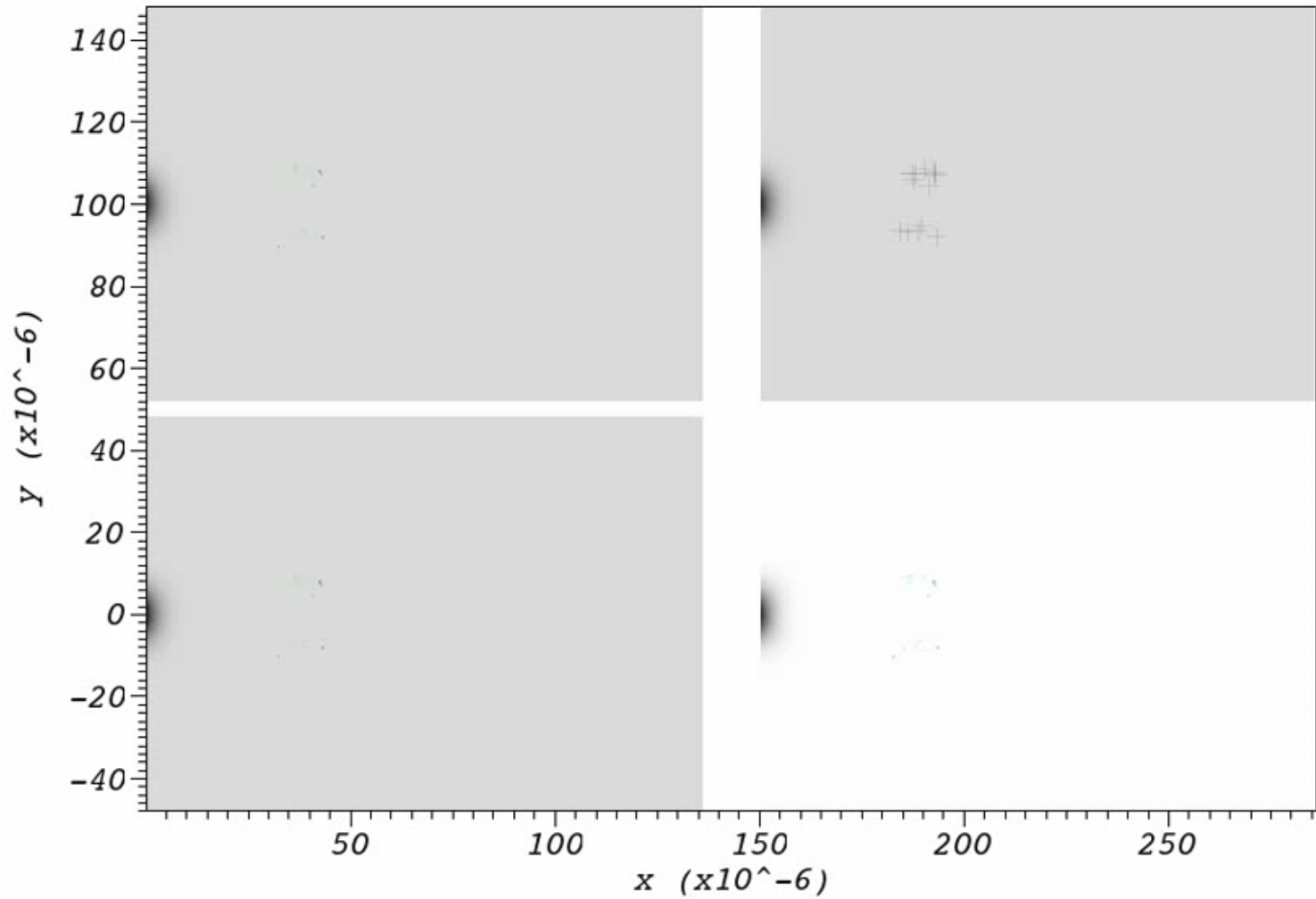
Multi-bunch production

Ion shield suggests focus position/release order & distance between rls positions $\zeta_{\text{rls},n} < \zeta_{\text{rls},n+1}$
in co-moving and lab frame

$$z_{\text{rls},1} < z_{\text{rls},2} < z_{\text{rls},3}$$

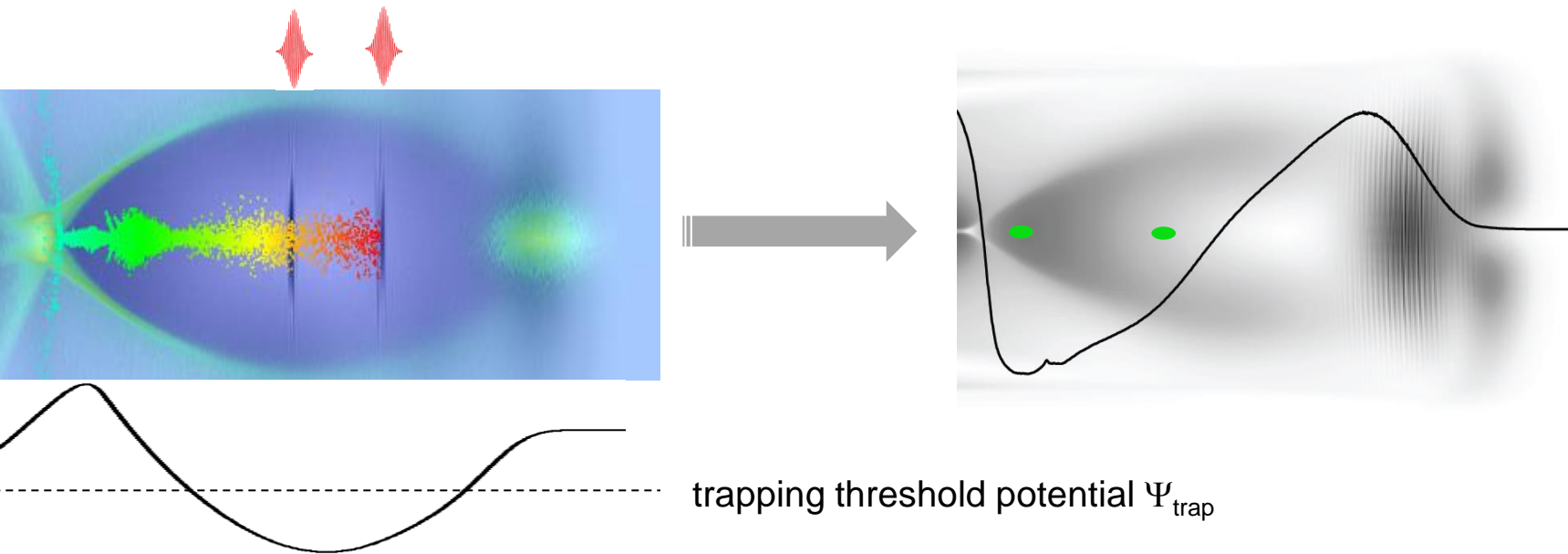


Photocathode + Space Charge Screening



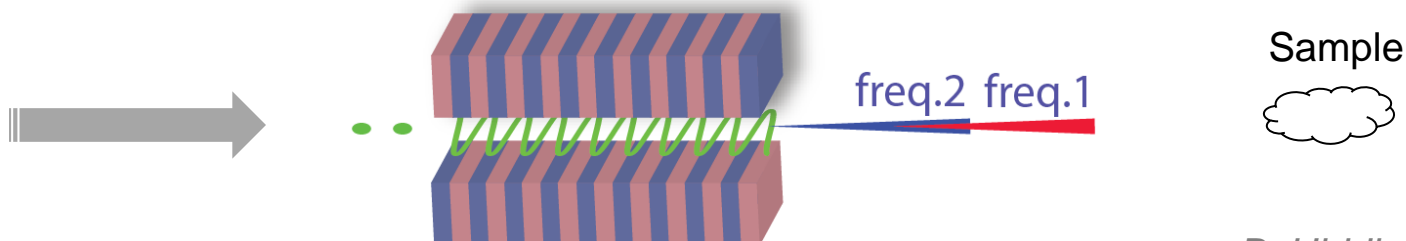
Tunable, lowest emittance multi-color FEL

Multiple laser pulses (~50 μJ each) generate multiple bunches of highest quality, separated by few μm



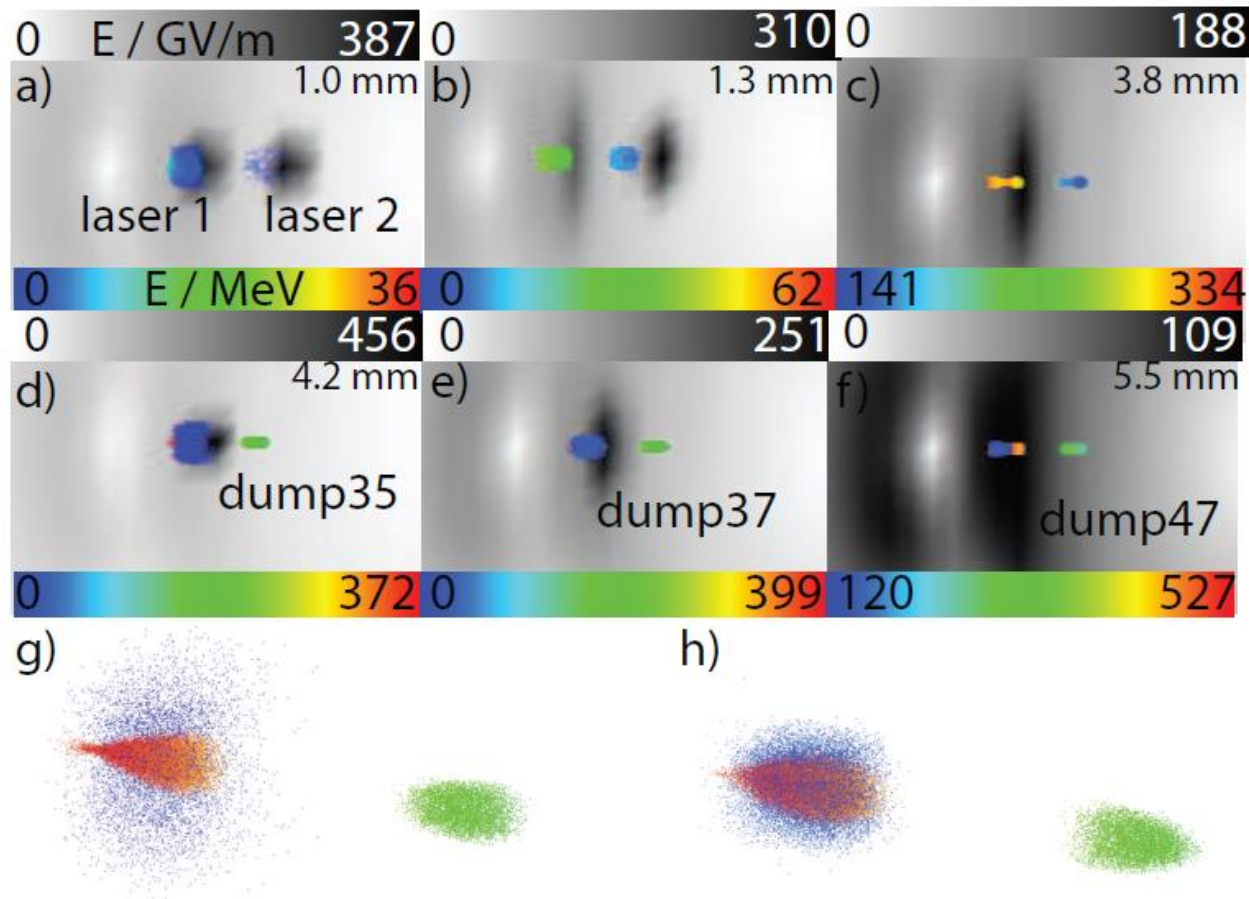
Electron bunch energies correlated to inter-bunch distance & release position, tunable in wide range

(cryogenic) undulator



Multi-bunch production

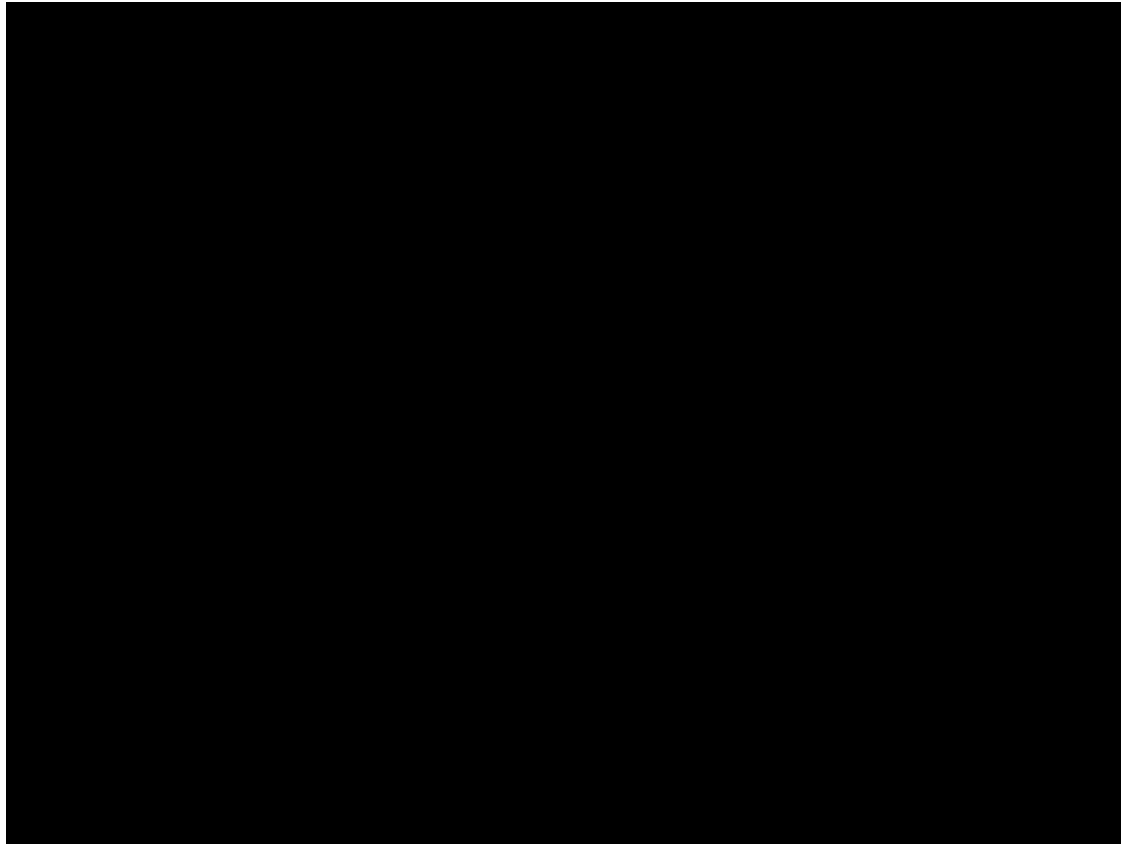
Energy tuning can be done by variation of release position in co-moving frame $\zeta = z - ct$ as well as in lab frame z



This way, fancy constellations can be produced: i.e. overlapping bunch hot and cold electron population, largely independent energy and delay tuning between bunches etc, bunch current shaping etc., tailored beam loading..

Preliminary PUFFIN 1D calculations
3d code: Phys. Plasmas 19, 093119 (2012)

B. McNeill, L. Campbell



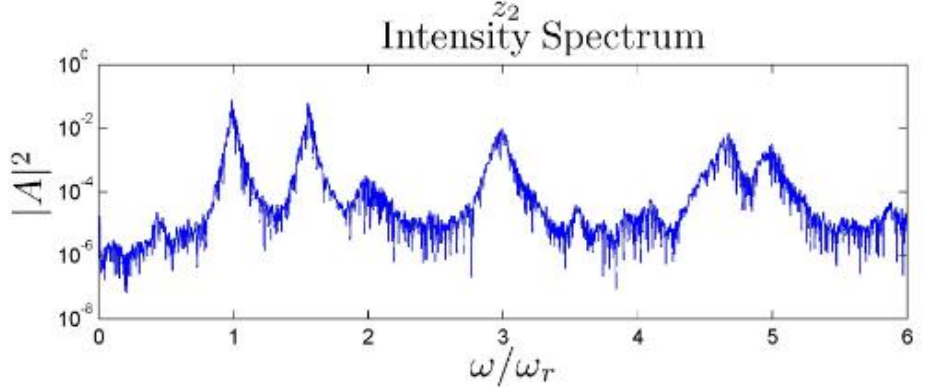
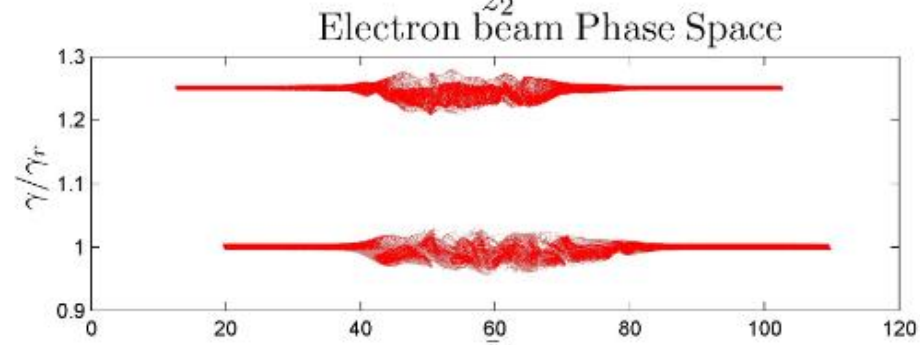
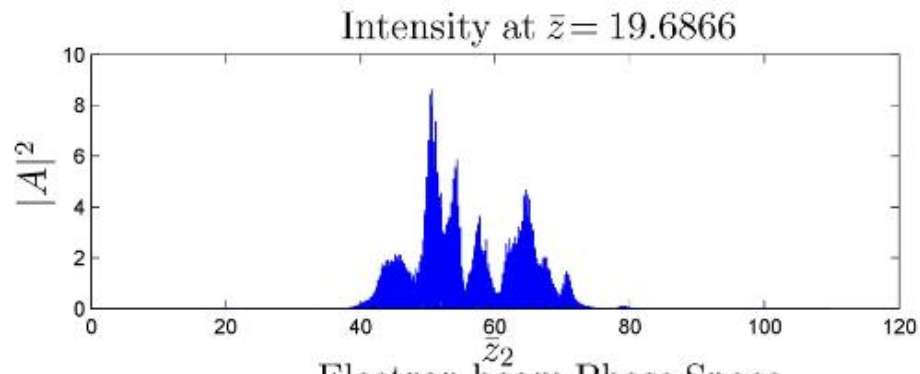
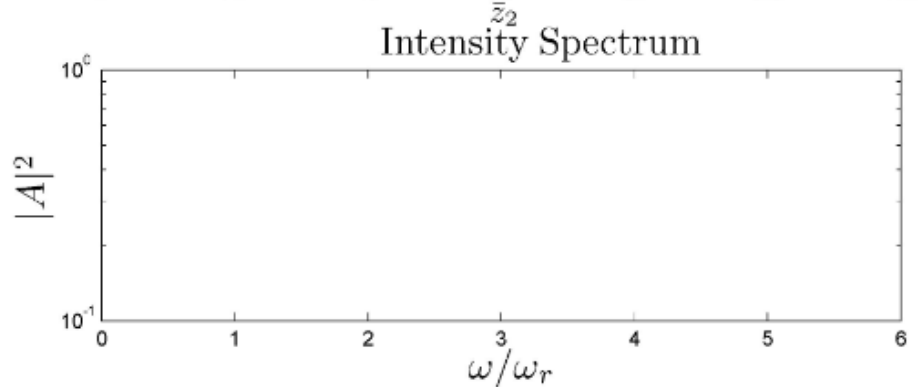
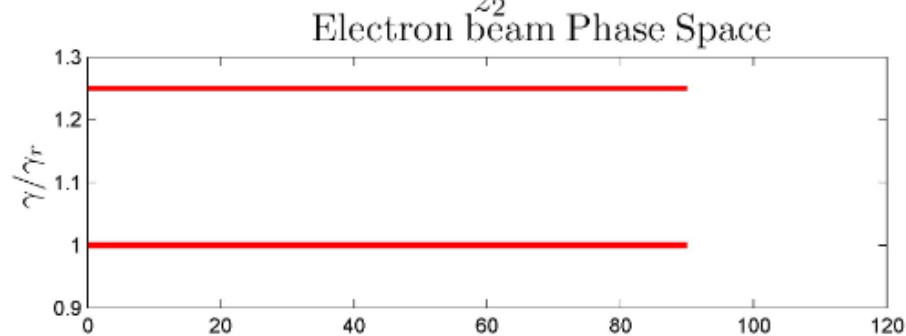
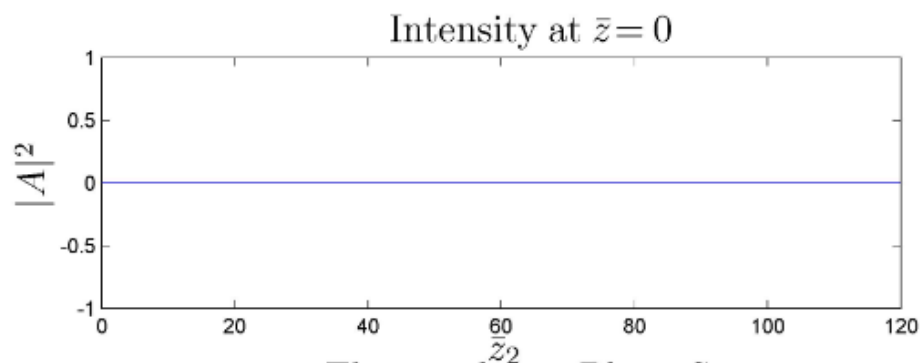
Preliminary PUFFIN 1D calculations

3d code: *Phys. Plasmas* 19, 093119 (2012)

B. McNeill, L. Campbell



Overlapping case:



Where/when to realize it?

FACET/SLAC

E-210 “Trojan Horse PWFA” expt., beamtime for 2014/15

- + stable driver beam
- + high energy beam
- + most extensive PWFA experience
- synchronization difficult
- Ionization/preionization difficult

Until recently!

Milestone: If this works, and if also the confidence level in extractable emittance is high enough → raise funding for Trojan FEL facilitie(s)

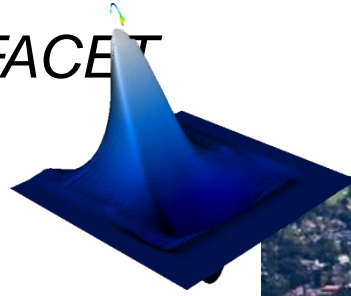
Photoinjector facilities (FLASHForward, FACET-II, SINBAD, CLARA?...)

- + very stable beam, high rep rate
- no facility online yet / no plasma acc. expmts. done yet
- not before 2016 (FLASHforward..)

Laser-Plasma-Accelerators worldwide (Strathclyde, Frascati, Jena, RAL, UHH/DESY...?)!

- + availability & cost-effectiveness
- + inherent perfect synchronization between electron bunch and Trojan release pulse
- instable performance
- no purposeful PWFA experiment has been demonstrated yet
- so far low (10 Hz) rep rate for 100 MeV+ beams

E-210 Trojan Horse PWFA @ FACET

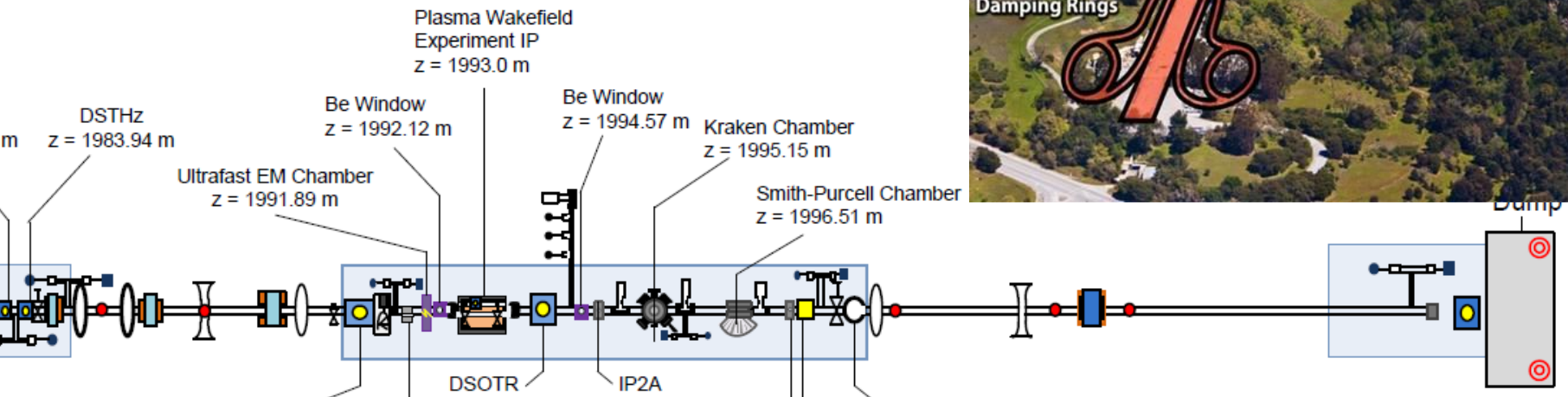
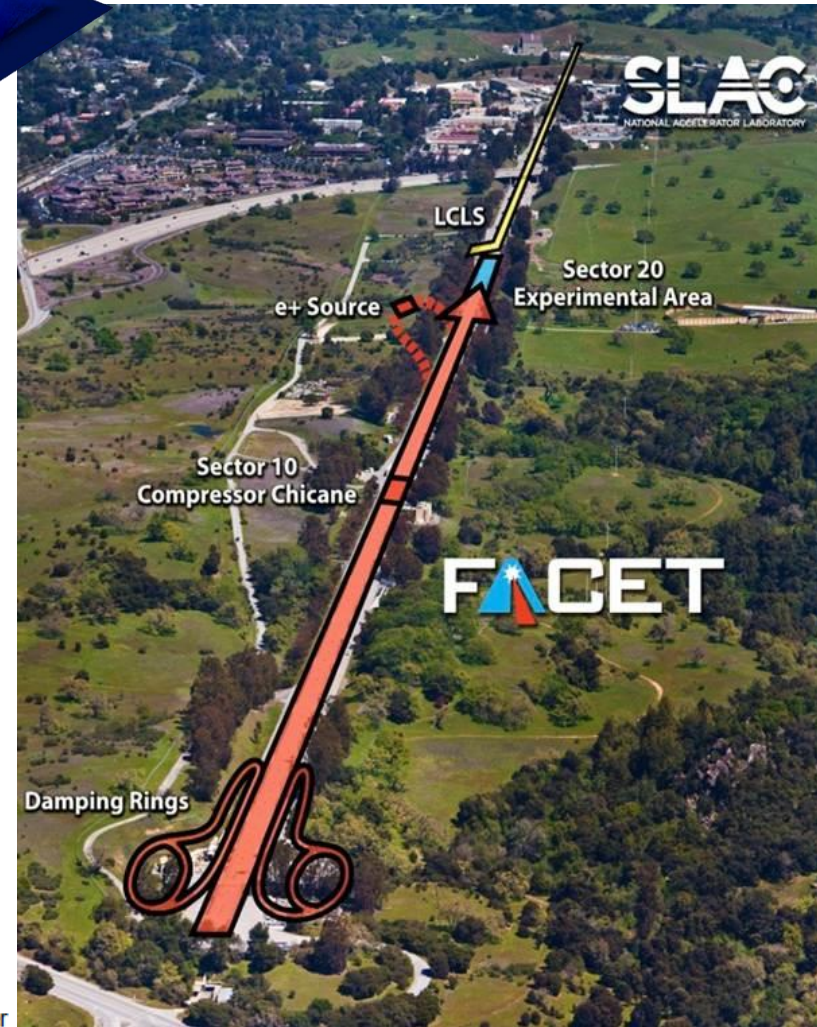


Pros:

- stable electron driver beam, can self-ionize Rb
- High energy bunch: 23 GeV
- 10-TW Laser system to be installed (until May 2013) for preionization (E-200 expt.) and E-210 expt.

Cons:

- Laser-jitter to master clock expected to be ± 40 fs
- Electron beam jitter < 1 ps
- Needs 2 km accelerator



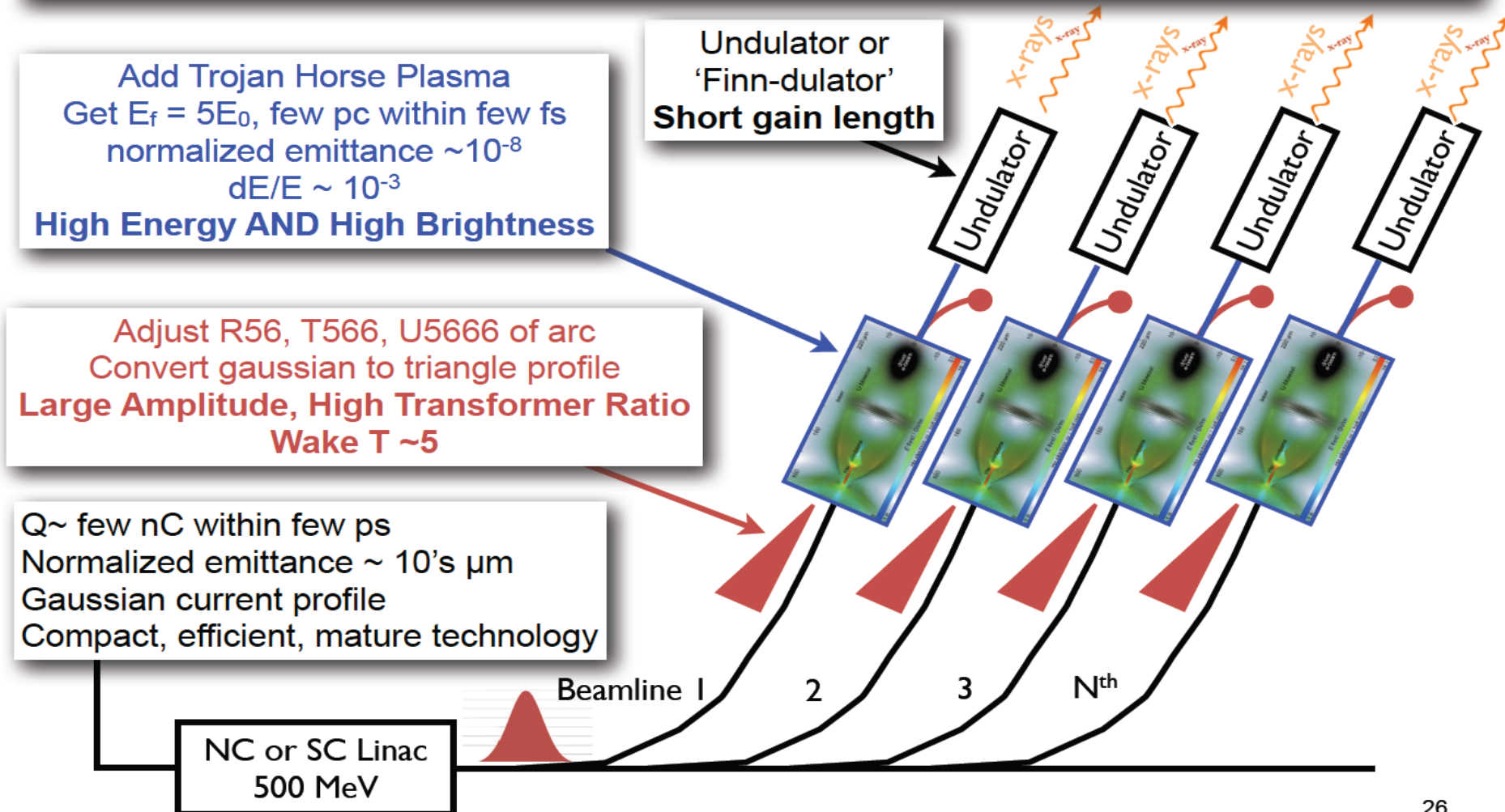
Hybrid Trojan Horse-based Future FEL facility?

w/ M. Hogan (SLAC) et al., 5th Generation Light Source Workshop, 2013

A Plasma Wakefield Accelerator Driven Compact X-FEL Plasma is Energy AND Brightness Transformer

SLAC

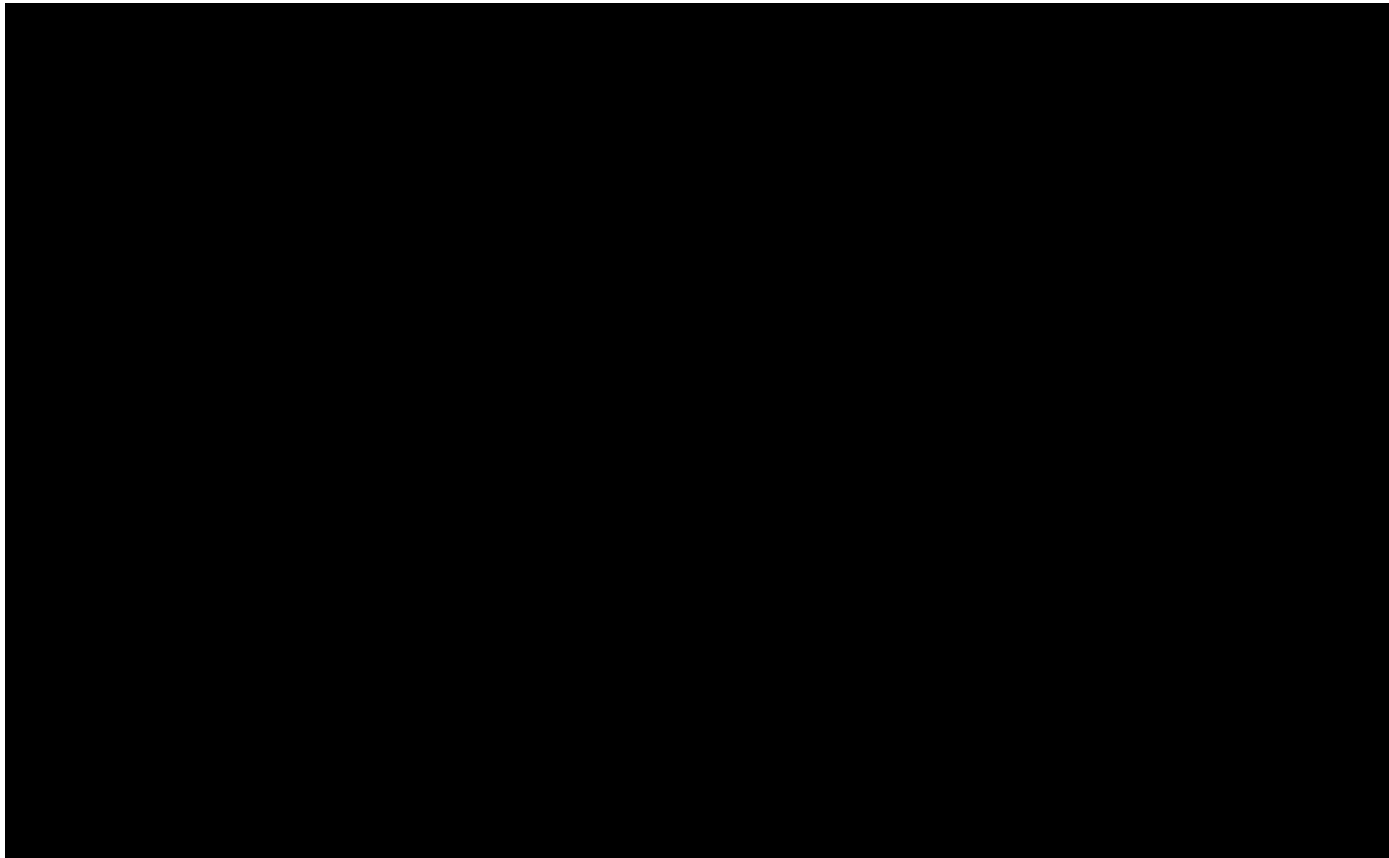
Compactness of plasmas accelerators, rep rate like rings with high brightness of linacs!



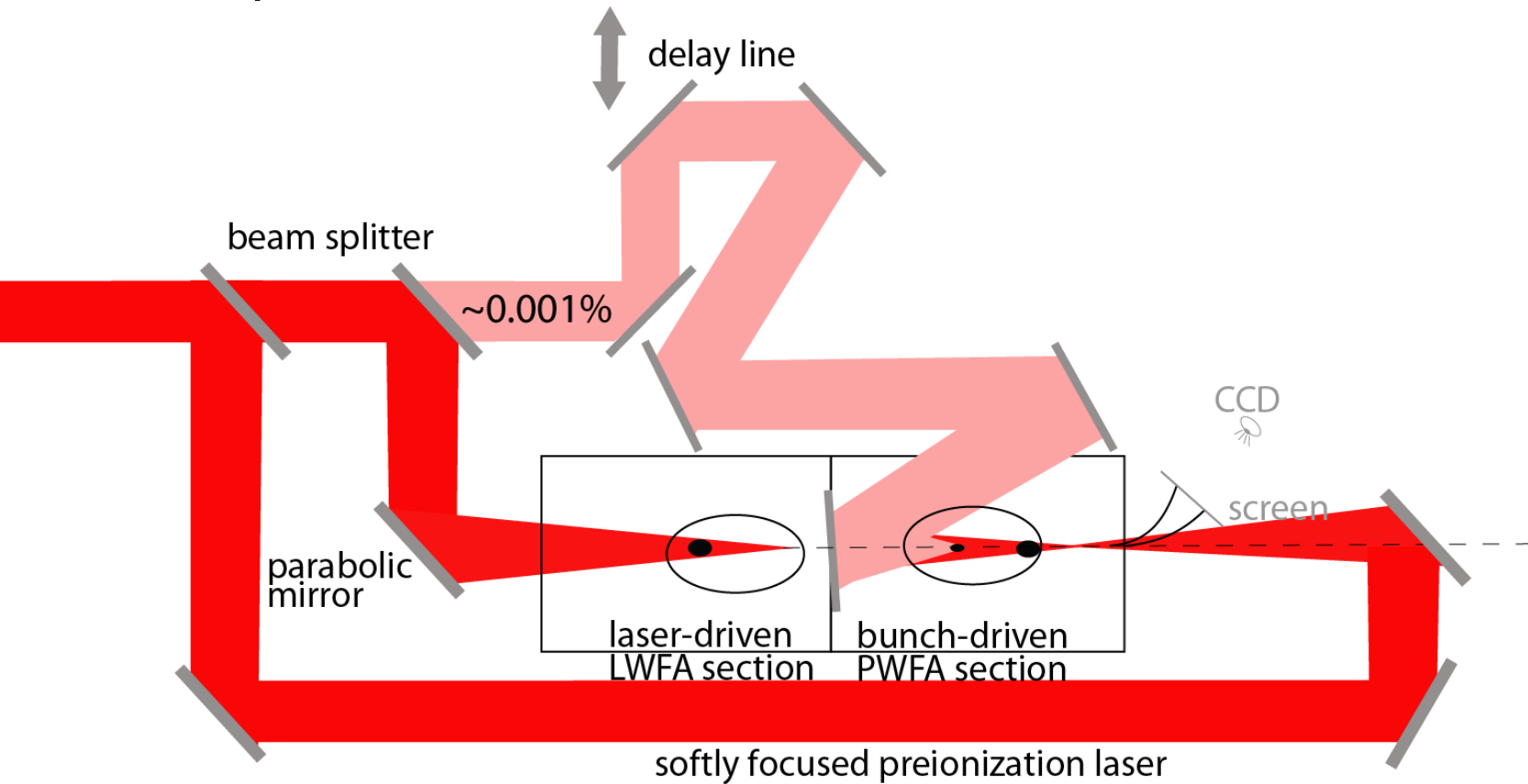
Scheme does also work w/ other high-brightness beam drivers, e.g. w/ FLASH driver:

Beam driver: $Q \approx 180$ pC, $\sigma_z \approx 8.39$ μm (rms), $\sigma_{x,y} \approx 7$ μm (rms), $E \approx 1.2$ GeV, $\Delta E \approx 0.1\%$,
Medium: LIT: preionized plasma $n_{\text{LIT}} \approx 1.5 \times 10^{17}$ cm^{-3} , HIT: He (IP 24.6 eV), $n_{\text{HIT}} \approx 3 \times 10^{17}$ cm^{-3}
Laser: $\tau \approx 25$ fs, $\lambda = 0.8$ μm , $w_0 \approx 5$ $\mu\text{m} \Rightarrow Z_R \approx 10$ μm , $a_0 \approx 0.03$, $I \approx 1.9 \times 10^{15}$ W/cm²

Experiments to start 2016



Beam brightness transformer and stabilizer for Laser-plasma-accelerators



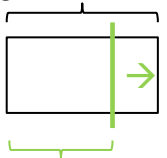
- Bunch quality transformer: energy, energy spread (see “Monoenergetic energy doubling”, PRL 140195002, 2010), emittance
- e.g., LPA: $\Delta E_1 = 20\%$, $\varepsilon_{n1} \sim 10^{-6}$ m rad \rightarrow TROJAN: $\Delta E_2 = 0.1\%$, $\varepsilon_{n2} \sim 10^{-8}$ m rad

Substantially different parameter goals for electron bunches from LWFA to be used for PWFA

- Energy spread does not matter, as long as the energy is sufficiently relativistic: all electrons move with $\sim c$. In first approximation, a 1 GeV bunch with perfect energy spread won't drive a different plasma wake than a 1 GeV bunch with 30% spread

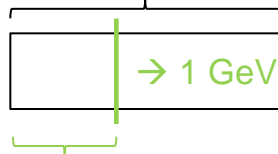
- Energy stability not so important: cap acceleration distance via preionization

max. possible acc.
length in this shot



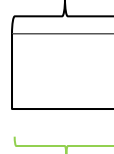
→ 1 GeV

max. possible acc.
length in this shot



→ 1 GeV

max. possible acc.
length in this shot



→ 1 GeV

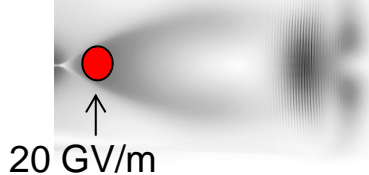
capped acc. length

capped acc. length

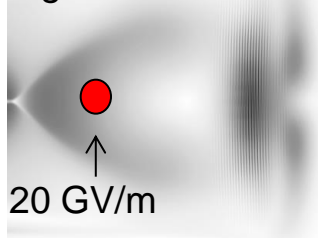
capped acc. length

- Even in case of current jitter, some stabilization is automatically achieved by the trapping process:

low driver current



high driver current



Even though the trapping position will be different, the acc. fields can be the same / very similar!

- Produce compact bunches with a lot of charge (via downramp injection?)
- Improve pointing
- Prevent dark current
- Can we somehow produce current upramp with LPA's for enhanced transformer ratio? ...

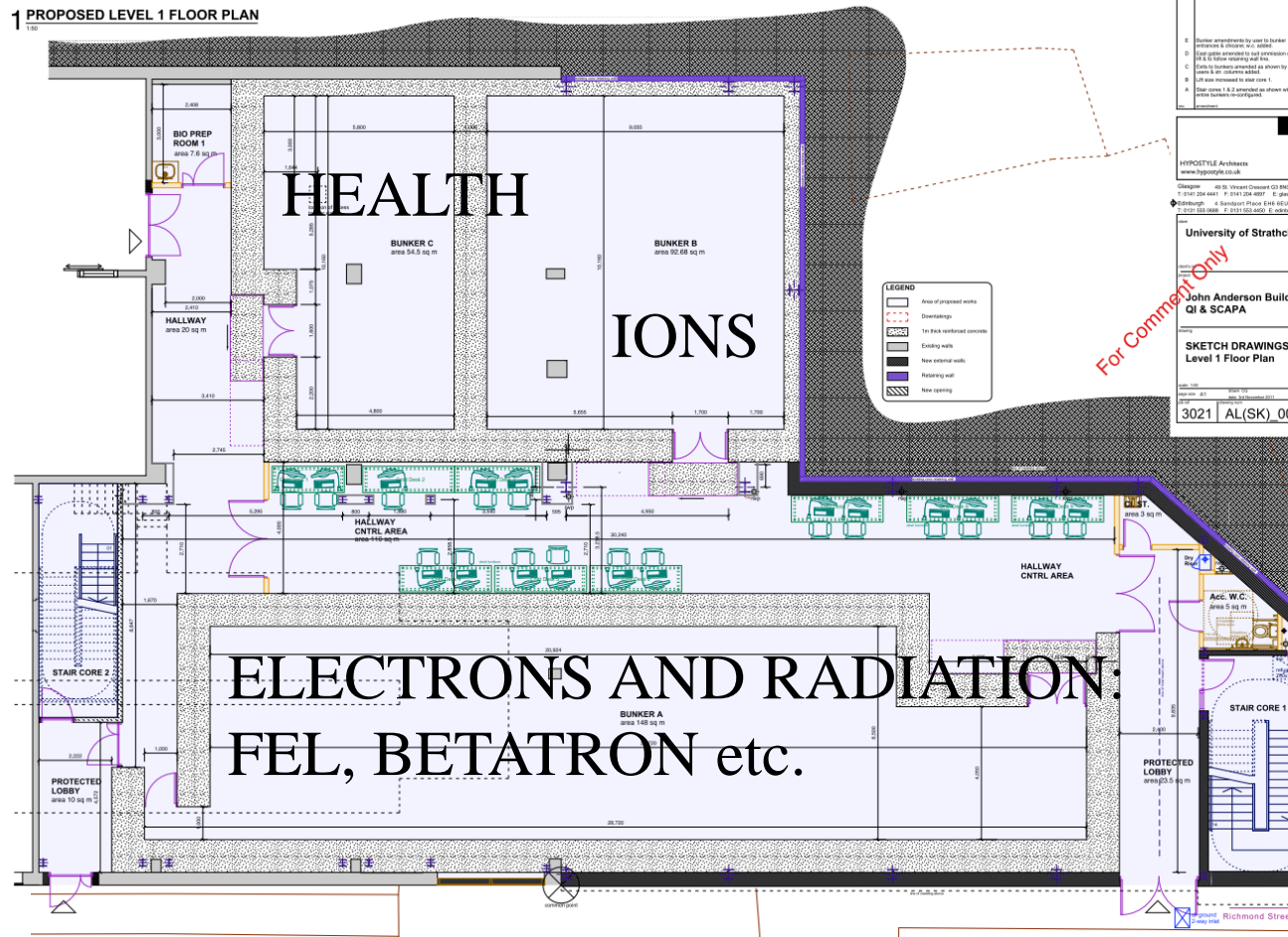


The Technology
& Innovation Centre

APPLICATIONS

- Radiobiology
- Ultrafast Probing
- High-Resolution Imaging
- Radioisotope Production
- Detector Development
- Radiation Damage Testing

1 PROPOSED LEVEL 1 FLOOR PLAN



APPLICATIONS

- Radiobiology
- Ultrafast Probing
- High-Resolution Imaging
- Radioisotope Production
- Detector Development
- Radiation Damage Testing

1200 m2 laboratory space, 3 shielded areas with 7 beam lines. 200-300 TW laser, 40 TW laser, sub-TW laser

Director: Dino Jaroszynski

Other key people: Paul McKenna, Zheng-Ming Sheng, Mark Wiggins, Gregory Welsh, Brian McNeill, Bernhard Hidding..

5th Generation Light Sources...

...need a 4th Generation Electron Source

5th
4D, high
wavelength,
compact?

4th
Free-Electron-
Laser

3rd
Undulator
radiation

2nd
Synchrotron
radiation

1st
Bremsstrahlung

4th
10's of GV/m fields in
plasmas & underdense
photocathode PWFA

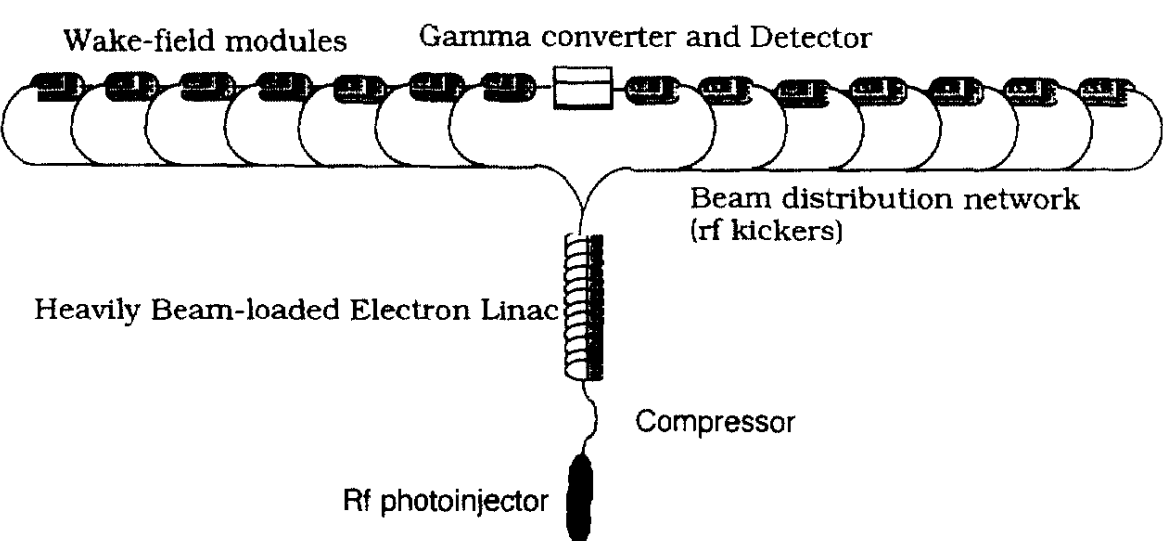
3rd
10's of GV/m fields in
plasmas
(LWFA and PWFA)

2nd
10's of MV/m fields,
photocathode
(e.g. FLASH, LCLS, XFEL)

1st
10's of MV/m fields,
thermionic cathode
(e.g. SLAC)

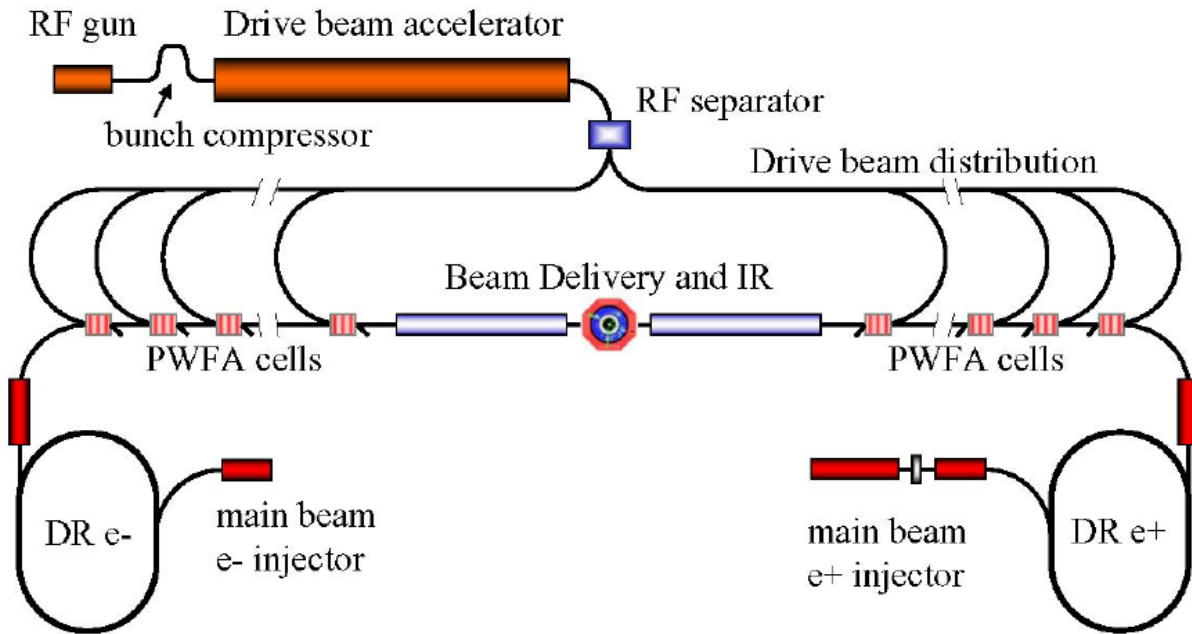
3rd way: All-optically powered TeV accelerator

PWFA-based:



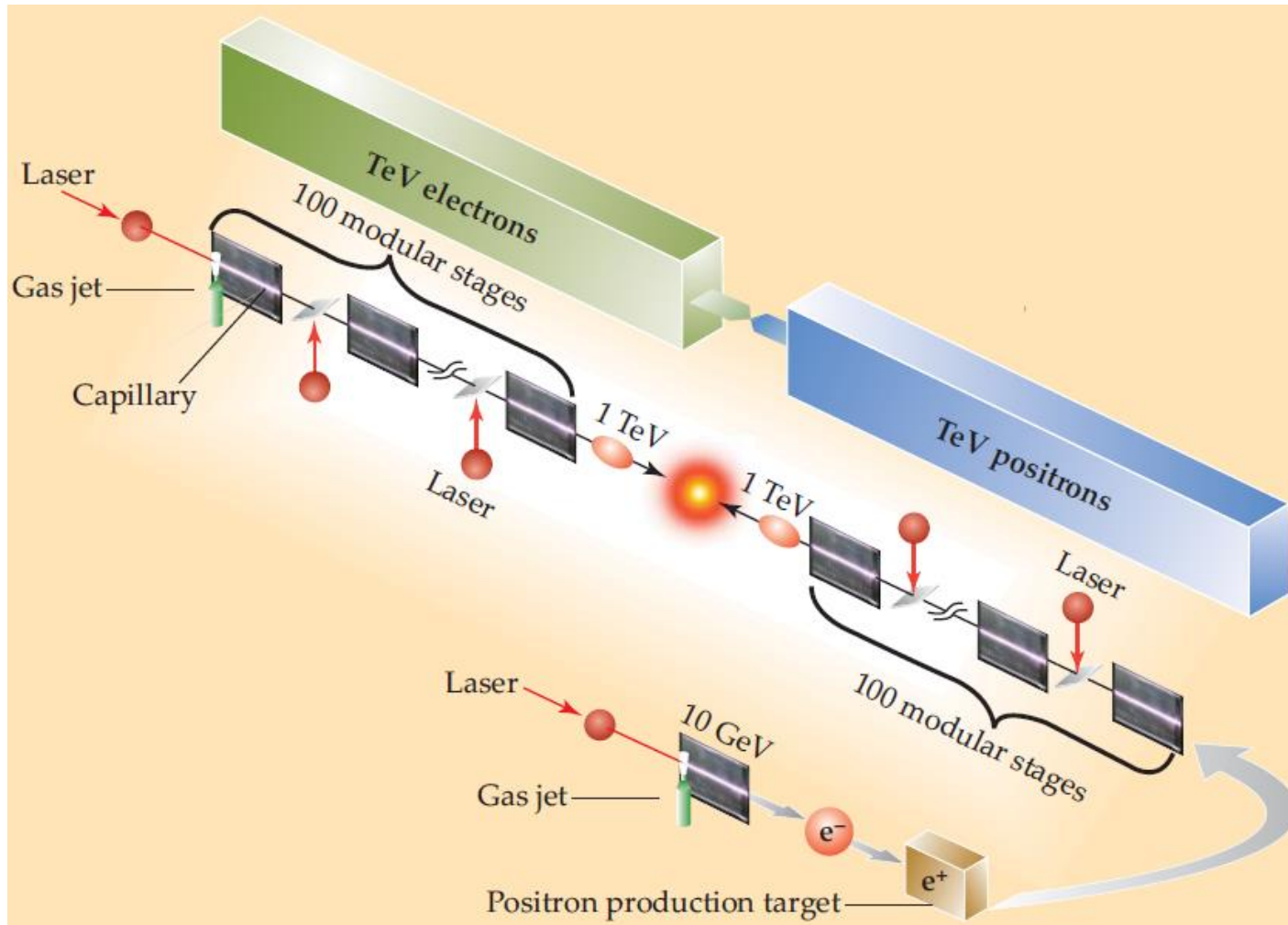
γ - γ - collider
Towards a Plasma Wake-field
Acceleration-based Linear Collider,
J.B. Rosenzweig et al.,
NIMA 410, 1998

e-e⁺ collider
A Concept of Plasma Wakefield
Acceleration Linear Collider (PWFA-LC),
A. Seryi et al.,
PAC 2009



3rd way: All-optically powered TeV accelerator

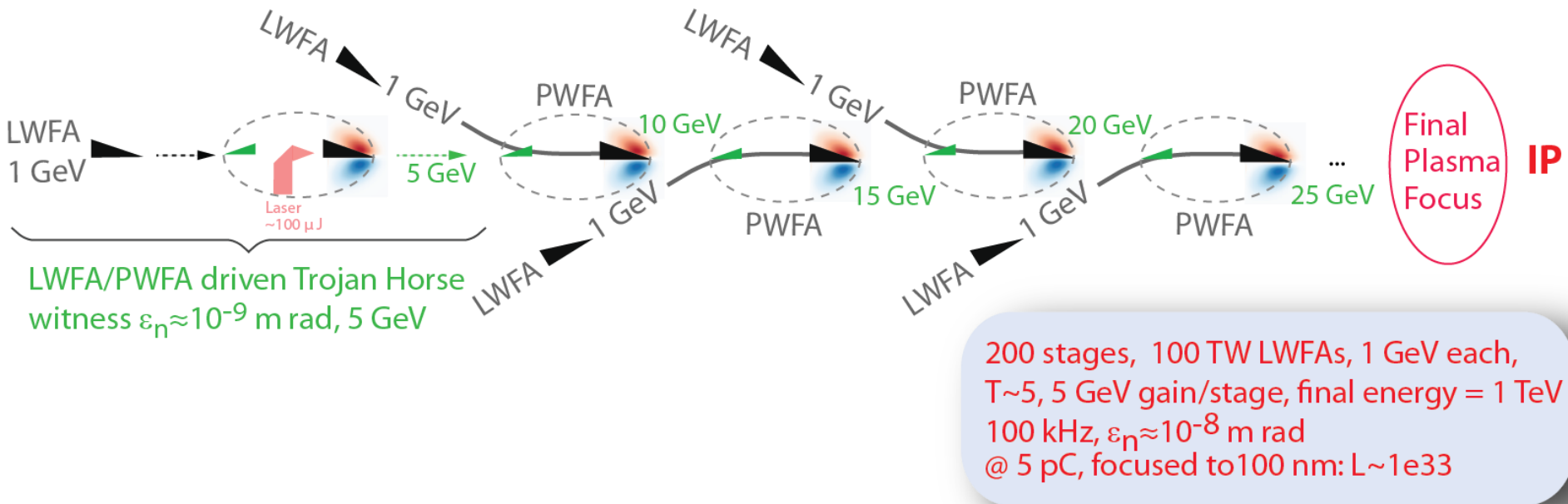
LWFA-based:



W. Leemans et al.,
Phys. Today, 2009

3rd way: All-optically powered TeV accelerator

Hybrid-LWFA/PWFA-based:



e⁻ e⁺ collider? $\gamma\gamma$ -collider?

..to be submitted

PCT International Patent Application

Title: METHOD FOR GENERATING ELECTRON BEAMS IN A
HYBRID LASER-PLASMA ACCELERATOR

Inventor(s): Bernhard Hidding et al.

Priority Date: June 18, 2011

Filing Date: June 18, 2012

Ser. No.: PCT/US12/043002

PRL 108, 035001 (2012)

PHYSICAL REVIEW LETTERS

week ending
20 JANUARY 2012

Ultracold Electron Bunch Generation via Plasma Photocathode Emission and Acceleration in a Beam-Driven Plasma Blowout

B. Hidding,^{1,2} G. Pretzler,² J. B. Rosenzweig,¹ T. Königstein,² D. Schiller,¹ and D. L. Bruhwiler³

¹*Department of Physics and Astronomy, University of California Los Angeles, Los Angeles, California 90095, USA*

²*Institut für Laser- und Plasmaphysik, Heinrich-Heine-Universität Düsseldorf, 40225 Düsseldorf, Germany*

³*Tech-X Corporation, Boulder, Colorado 80303, USA*

(Received 30 March 2011; published 17 January 2012)

Beam-driven plasma wakefield acceleration using low-ionization-threshold gas such as Li is combined with laser-controlled electron injection via ionization of high-ionization-threshold gas such as He. The He electrons are released with low transverse momentum in the focus of the copropagating, nonrelativistic-intensity laser pulse directly inside the accelerating or focusing phase of the Li blowout. This concept paves the way for the generation of sub- μm -size, ultralow-emittance, highly tunable electron bunches, thus enabling a flexible new class of an advanced free electron laser capable high-field accelerator.

Beyond Injection: Trojan Horse Underdense Photocathode Plasma Wakefield Acceleration

B. Hidding^{*,†}, J.B. Rosenzweig[†], Y. Xi[†], B. O'Shea[†], G. Andonian[†], D. Schiller[†], S. Barber[†], O. Williams[†], G. Pretzler^{*}, T. Königstein^{*}, F. Kleeschulte^{*}, M. J. Hogan^{**}, M. Litos^{**}, S. Corde^{**}, W. W. White^{**}, P. Muggli[‡], D.L. Bruhwiler^{§,¶} and K. Lotov^{||,††}

**Institut für Laser- und Plasmaphysik, Heinrich-Heine-Universität Düsseldorf, 40225 Düsseldorf, Germany*

†Particle Beam Physics Laboratory, Department for Physics and Astronomy, UCLA, USA

Citation: [AIP Conf. Proc. 1507, 570 \(2012\)](#); doi: [10.1063/1.4773760](#)

View online: <http://dx.doi.org/10.1063/1.4773760>

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|| Budker Institute of Nuclear Physics SB RAS, 630090, Novosibirsk, Russia

†† Novosibirsk State University, 630090, Novosibirsk, Russia

Abstract. An overview on the underlying principles of the hybrid plasma wakefield acceleration scheme dubbed "Trojan Horse" acceleration is given. The concept is based on laser-controlled release of electrons directly into a particle-beam-driven plasma blowout, paving the way for controlled, shapeable electron bunches with ultralow emittance and ultrahigh brightness. Combining the virtues of a low-ionization-threshold underdense photocathode with the GV/m-scale electric fields of a practically dephasing-free beam-driven plasma blowout, this constitutes a 4th generation electron acceleration scheme. It is applicable as a beam brightness transformer for electron bunches from LWFA and PWFA systems alike. At FACET, the proof-of-concept experiment "E-210: Trojan Horse Plasma Wakefield Acceleration" has recently been approved and is in preparation. At the same time, various LWFA facilities are currently considered to host experiments aiming at stabilizing and boosting the electron bunch output quality via a trojan horse afterburner stage. Since normalized emittance and brightness can be improved by many orders of magnitude, the scheme is an ideal candidate for light sources such as free-electron-lasers and those based on Thomson scattering and betatron radiation alike.

Hybrid modeling of relativistic underdense plasma photocathode injectors

Y. Xi,¹ B. Hidding,^{1,2} D. Bruhwiler,³ G. Pretzler,⁴ and J. B. Rosenzweig¹

¹*Department of Physics and Astronomy, University of California, Los Angeles, California, USA*

²*Institut für Experimentalphysik, Universität Hamburg & DESY, 22607 Hamburg, Germany*

³*University of Colorado at Boulder, 390UCB, Boulder, Colorado 80309, USA*

⁴*Institut für Laser- und Plasmaphysik, Heinrich-Heine-Universität Düsseldorf, Germany*

(Received 6 November 2012; published 25 March 2013)

The dynamics of laser ionization-based electron injection in the recently introduced plasma photocathode concept is analyzed analytically and with particle-in-cell simulations. The influence of the initial few-cycle laser pulse that liberates electrons through background gas ionization in a plasma wakefield accelerator on the final electron phase space is described through the use of Ammosov-Delone-Krainov theory as well as nonadiabatic Yudin-Ivanov (YI) ionization theory and subsequent downstream dynamics in the combined laser and plasma wave fields. The photoelectrons are tracked by solving their relativistic equations of motion. They experience the analytically described transient laser field and the simulation-derived plasma wakefields. It is shown that the minimum normalized emittance of fs-scale electron bunches released in multi-GV/m-scale plasma wakefields is of the order of 10^{-2} mm mrad. Such unprecedented values, combined with the dramatically increased controllability of electron bunch production, pave the way for highly compact yet ultrahigh quality plasma-based electron accelerators and light source applications.

Low transverse emittance electron bunches from two-color laser-ionization injection

Lu-Le Yu^{ab}, Eric Esarey^a, Jean-Luc Vay^a, Carl B. Schroeder^a, Carlo Benedetti^a, Cameron G. R. Geddes^a, Sergey G. Rykovanov^a, Stepan S. Bulanov^b, Min Chen^c and Wim P. Leemans^{ab}

^aLawrence Berkeley National Laboratory, Berkeley, California 94720, USA;

^bUniversity of California, Berkeley, California 94720, USA;

^cShanghai Jiao Tong Univeristy, Shanghai 200240, China

ABSTRACT

A method is proposed to generate low emittance electron bunches from two color laser pulses in a laser-plasma accelerator. A two-region gas structure is used, containing a short region of a high-Z gas (e.g., krypton) for ionization injection, followed by a longer region of a low-Z gas for post-acceleration. A long-laser-wavelength (e.g., 5 μm) pump pulse excites plasma wake without triggering the inner-shell electron ionization of the high-Z gas due to low electric fields. A short-laser-wavelength (e.g., 0.4 μm) injection pulse, located at a trapping phase of the wake, ionizes the inner-shell electrons of the high-Z gas, resulting in ionization-induced trapping. Compared with a single-pulse ionization injection, this scheme offers an order of magnitude smaller residual transverse momentum of the electron bunch, which is a result of the smaller vector potential amplitude of the injection pulse.

Phase Space Dynamics of Ionization Injection in Plasma Based Accelerators

X. L. Xu,¹ J. F. Hua,¹ F. Li,¹ C. J. Zhang,¹ L. X. Yan,¹ Y. C. Du,¹ W. H. Huang,¹
H. B. Chen,¹ C. X. Tang,¹ W. Lu,^{1,2,*} P. Yu,² W. An,² W. B. Mori,² and C. Joshi²

¹Key Laboratory of Particle and Radiation Imaging of Ministry of Education, Tsinghua University, Beijing 100084, China

²University of California Los Angeles, LA, CA 90095, USA

(Dated: June 19, 2013)

The evolution of beam phase space in ionization-induced injection into plasma wakefields is studied using theory and particle-in-cell (PIC) simulations. The injection process causes special longitudinal and transverse phase mixing leading initially to a rapid emittance growth followed by oscillation, decay, and eventual slow growth to saturation. An analytic theory for this evolution is presented that includes the effects of injection distance (time), acceleration distance, wakefield structure, and nonlinear space charge forces. Formulas for the emittance in the low and high space charge regimes are presented. The theory is verified through PIC simulations and a good agreement is obtained. This work shows how ultra-low emittance beams can be produced using ionization-induced injection.

Two-Pulse Ionization Injection into Quasi-Linear Laser Wakefields

N. Bourgeois, J. Cowley, and S. M. Hooker

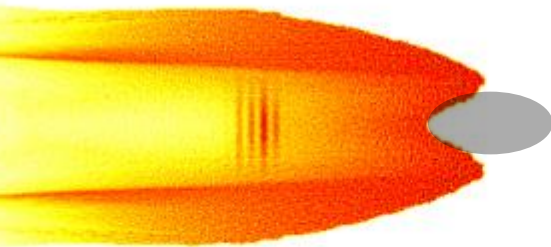
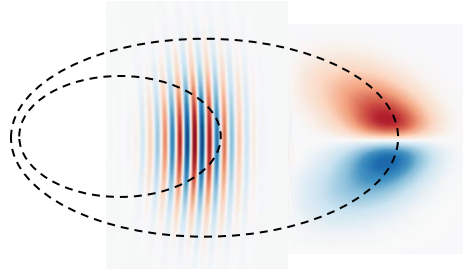
*Department of Physics, University of Oxford, Clarendon Laboratory,
Parks Road, Oxford OX1 3PU, United Kingdom*

(Dated: July 1, 2013)

We describe a scheme for controlling electron injection into the quasi-linear wakefield driven by a guided drive pulse via ionization of a dopant species by a collinear injection laser pulse with a short Rayleigh range. The scheme is analyzed by particle in cell simulations which show controlled injection and acceleration of electrons to an energy of 370 MeV, a relative energy spread of 2%, and a normalized transverse emittance of $3.0 \mu\text{m}$.

Trojan Horse (Pre)History

2008: Laser-driven bubble in a beam-driven blowout?
("Matryoshka acc..")



Laser pulse at typical (relativistic) LWFA intensities expel electrons



2008/2009: much better mode would be to have the laser pulse at minimal intensity ($a_0 \ll 1$), so that released electrons are "still" and remain still inside the blowout \rightarrow "Trojan horse acc.", originally to be presented at AAC 2010 in Kardamili, Greece (sic!)



2010: spin -off idea: PWFA with electron beams from LWFA ("Hybrid energy doubling", PRL 104, 195002, 2010)

2011: DE patent, PRL submitted..

Summary

- electron bunches with unprecedented emittance ($\varepsilon_n \sim 10^{-9} - 10^{-10}$ m rad) and brightness may be possible (emittance preservation & extraction crucial)
- Unprecedented bunch shaping capabilities (more flexible than state-of-the-art photoinjectors)
- Trojan horse a bunch quality transformer (e.g., $\Delta E_1 = 20\%$, $\varepsilon_{n1} \sim 10^{-6}$ m rad $\rightarrow \Delta E_2 = 20\%$, $\varepsilon_{n2} \sim 10^{-8}$ m rad)
- ... and as a bunch energy transformer
- Output stabilizer for LWFA
- Scheme applicable for most diverse scenarios: hybrid conventional/PWFA accelerator (SLAC/Trojan Horse), hybrid photoinjector/PWFA accelerator (FLASH, FACET-II, CLARA etc.), hybrid LWFA/PWFA accelerators
- FEL game changer? Performance **substantially** better than XFEL may be possible:
Reduce costs from ~1 Mrd. € to ~5 Mio. €, size from km to m-scale, yet better performance (e.g., multi-color FELs)
- HEP accelerator applications? TeV accelerators..

Collaborators:



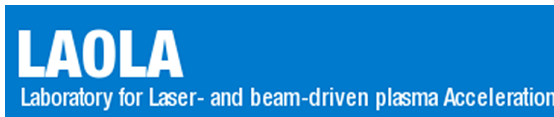
D. Jaroszynski, ZM Sheng, P. McKenna, B. McNeill, L. Campbell, M. Wiggins, G. Welsh, G. Manahan, G. McKendrick et al.



J.B. Rosenzweig, D. Bruhwiler, Y. Xi, A. Deng, G. Andonian, D. Schiller, B. O'Shea, S. Barber, O. Williams et al.



O. Karger, C. Aniculaesei, G. Wittig, T. Heinemann, G. Fuhs, J. Wein, M. Quast, H. Groth, T. Kovener, F. Habib, P. Scherkl, G. Hurtig



laola.desy.de

